

# Modelling the Electrical Activity of Pancreatic Beta Cells

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*A dissertation submitted for the partial fulfilment of BS-MS Dual  
Degree in Science*



Indian Institute of Science Education and Research, Mohali  
April 2016

## Certificate of Examination

This is to certify that the dissertation titled “**Modelling the Electrical Activity of Pancreatic Beta Cells**” submitted by **Ms. Yosman BapatDhar** (Reg. No. MS11060) for the partial fulfillment of BS-MS dual degree programme of the Institute, has been examined by the thesis committee duly appointed by the Institute. The committee finds the work done by the candidate satisfactory and recommends that the report be accepted.

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## Declaration

The work presented in this dissertation has been carried out by me under the guidance of Prof. Somdatta Sinha at the Indian Institute of Science Education and Research Mohali.

This work has not been submitted in part or in full for a degree, a diploma, or a fellowship to any other university or institute. Whenever contributions of others are involved, every effort is made to indicate this clearly, with due acknowledgment of collaborative research and discussions. This thesis is a bona-fide record of original work done by me and all sources listed within have been detailed in the bibliography.

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In my capacity as the supervisor of the candidate's project work, I certify that the above statements by the candidate are true to the best of my knowledge.

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## **Acknowledgements**

My sincere gratitude to Prof. Somdatta Sinha, my research supervisor, for her patient guidance, useful critiques of the work and assistance in keeping my progress on schedule.

I thank all the members of the Computational and Mathematical Biology Group for their inputs during discussion sessions. I would like to specially acknowledge Vivek Sagar's help in troubleshooting some of my programs.

Finally, a huge thank you to my family and all my friends for keeping me sane and smiling throughout the year.

## **Abstract**

The aim of this thesis project was to study the evolution of single cell mathematical models of the electrical activity of pancreatic beta cells, find a candidate model that best matched experimental data present in the existing literature, and to use it to model gap-junction coupled beta cells present in pancreatic islets. This study found that the Dual Oscillator Model is currently the best model for the electrical activity of a single beta cell. It was found that the behaviour of a 1D ring of cells with nearest neighbour coupling depends upon the number of cells in the ring, the percentage composition of each of the cell types and the coupling strength ( $g_c$ ) between neighbouring cells. Also, when multiple cells are coupled, synchronization of the burst goes from  $180^\circ$  anti phase (low  $g_c$ ) to out of phase (intermediate  $g_c$ ) to in phase (high  $g_c$ ).

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# Chapter 1

## Introduction

### 1.1 Electrically Active Cells

Electrically active cells are those cells that respond to changes in their membrane potential. Typically, their primary functions are regulated by changes in their membrane potential which occurs due to regulation of ionic fluxes. Examples are – neurons and conduction of action potentials; the electrical activity of cardiac cells for proper beating of the heart. Pancreatic beta cells are another example of electrically active cells. These cells respond to changes in their membrane voltage by secreting insulin.

#### 1.1.1 Pancreatic Islets of Langerhans

Beta cells are specialized cells present in large numbers in structures called *Islets of Langerhans* (Fig. 1.1) in the Pancreas. Each islet consists of thousands of cells. In an islet thousands of beta cells generally form the core, and other cells (delta and alpha cells secreting glucagon and somatostatin hormones respectively) are present in the periphery (Fig. 1.1).

Beta cells respond to elevated blood glucose levels by secreting the hormone Insulin in a pulsatile manner<sup>1,2</sup>. Coordinated membrane voltage bursting in beta cells

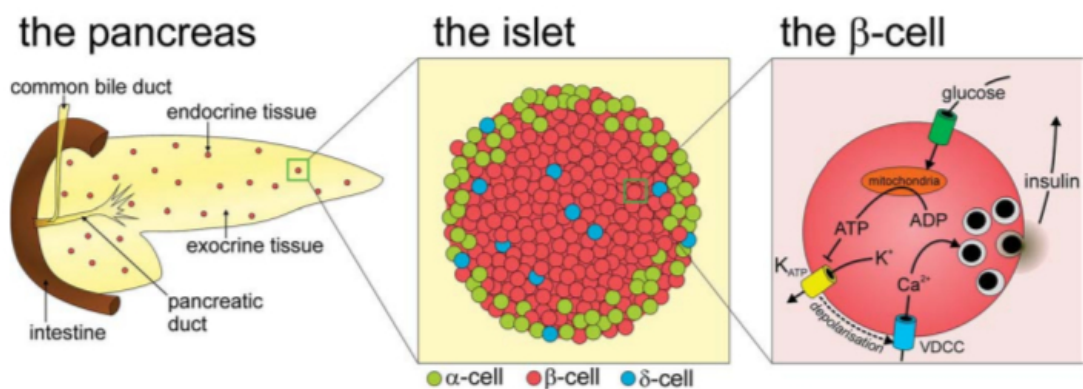


Figure 1.1: *The Pancreas, an Islet and a Beta Cell* [9]

(in an islet) (*Fig. 1.2*) precedes the rhythmic insulin secretion, and malfunction in the electrical activity of beta cells can lead to improper insulin secretion - a cause of Type 2 Diabetes<sup>3</sup>. Beta cells communicate (electrically) through membrane gap junctions (Connexin proteins). The major gap junction protein found between beta cells is Connexin-36<sup>6,7,8</sup>.

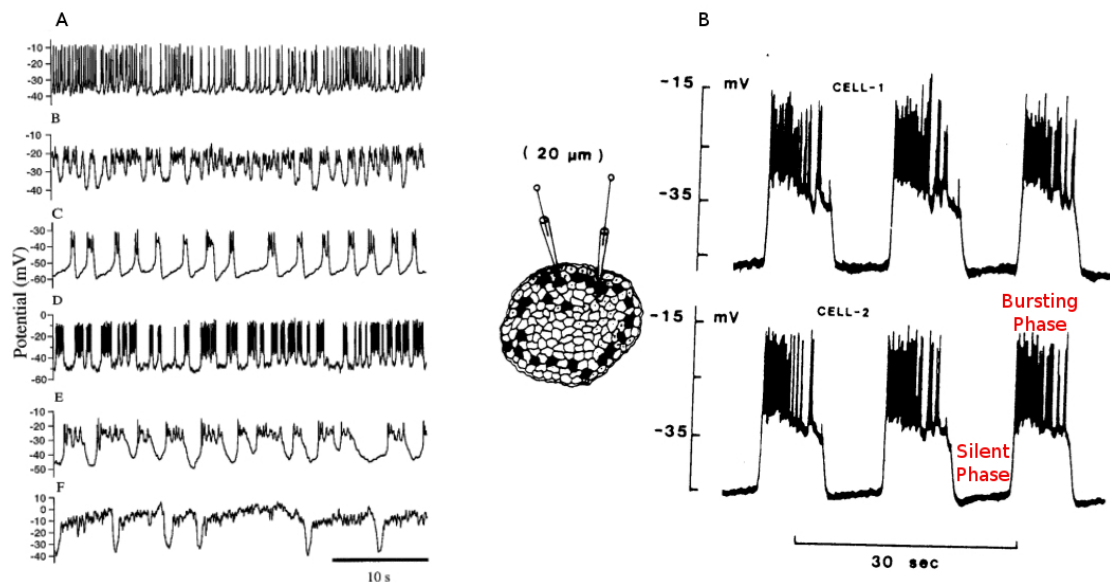


Figure 1.2: (A) Voltage time curves for single beta cells (B) Voltage time curves of 2 coupled cells in the islet [11]

Beta cell bursting consists of two phases: (1) A 'Silent' phase where the membrane voltage is approximately equal to the resting membrane voltage and (2) a 'Bursting' or an 'Active' phase where the membrane voltage fluctuates up and down rapidly (*Fig. 1.2 (B)*). The primary ionic currents responsible for membrane potential changes in beta cells are calcium and potassium currents of different types. Due to its involvement in diabetes, electrophysiological and pharmacological experiments and mathematical modelling of beta cell dynamics have been active areas of research.

## 1.2 Modelling the Biological Cell as an Electrical Circuit

The membrane of a biological cell provides a boundary separating the external environment from the internal environment of the cell (*Fig. 1.3(a)*). Typically, the membrane is selectively permeable thus regulating the passage of materials in and out of the cell. The cell membrane consists of a lipid bilayer (water insoluble) through which, irregularly interspersed, are globular proteins. The membrane also consists of protein lined water pores - channel proteins. Channel proteins (like voltage dependent calcium channels, VDCCs, ATP dependent potassium channels,  $K_{ATP}$  etc.) allow the passage of specific kinds of molecules. The selective permeability of the cell membrane and the presence of channel proteins results in a concentration difference across the membrane of various ions ( $Na^+$ ,  $K^+$ ,  $Ca^{+2}$ ,  $Cl^-$ ). These concentration differences are set-up and maintained by active processes that use ATP to pump ions against their electrochemical gradients.

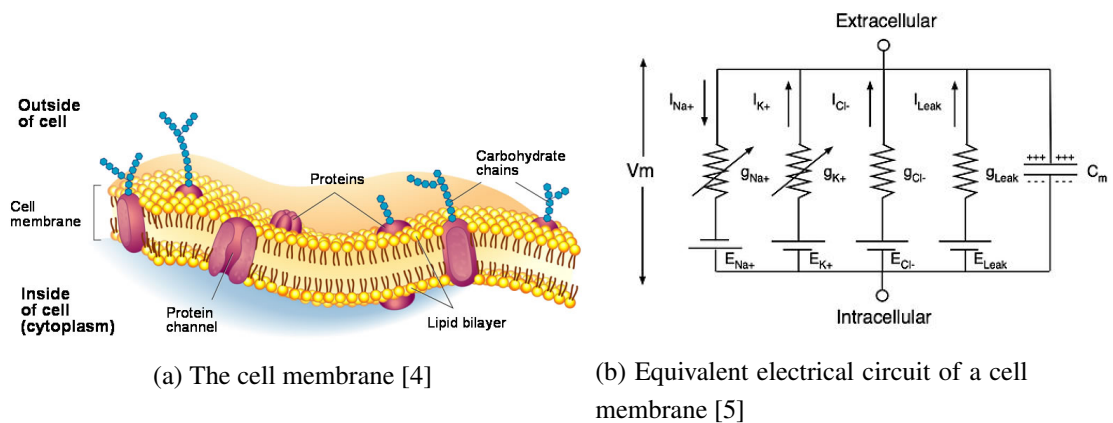


Figure 1.3: *The biological cell membrane as an electrical circuit*

Since the cell membrane separates and maintains a charge difference, it can be thought of as a capacitor (*Fig. 1.3(b)*). By the definition of the capacitance of a capacitor (ratio of charge across the capacitor to the voltage necessary to hold that charge), one can write:  $C_m = \frac{Q}{V}$

Since current is defined as  $\frac{dQ}{dt}$ , it follows that the capacitive current is (assuming  $C_m$  to be constant):

$$C_m \frac{dV}{dt} \quad (1.1)$$

For a given ionic species, say  $Na^+$ , the potential drop across the cell membrane has two components: first due to concentration differences, given by the Nernst equation ( $V_{Na} = \frac{RT}{ZF} \ln\left(\frac{[Na^+]_{external}}{[Na^+]_{internal}}\right)$ ) and, second, due to an electrical current  $rI_{Na}$  (if the channel is Ohmic), where  $r$  is the resistance of the channel. Summing the two contributions

we get:

$$V = rI_{Na} + V_{Na} \quad (1.2)$$

Solving for current we get:  $I_{Na} = g_{Na}(V - V_{Na})$  where  $g_{Na} = \frac{1}{r}$  is the channel conductance.

Since there is no net build-up of charge on either side of a cell's membrane, the sum of the ionic and capacitive charges should be zero. Hence:

$$C_m \frac{dV}{dt} + I_{ionic} = 0 \quad (1.3)$$

### 1.2.1 The Hodgkin-Huxley Model (1952)

In 1952, Alan Hodgkin and Andrew Huxley published their Nobel prize winning work as a final paper<sup>11</sup> in a series of papers. This last paper is a landmark in the field of mathematical physiology. In this paper, Hodgkin and Huxley (HH) present elegant experimental data, a theoretical hypothesis, a fit of the model to the data and, most crucially, a testable prediction. This work stands out as a rare example of a successful combination of experiment and theory!

HH developed the first model of the propagation of an electrical signal along the axon of a nerve in the giant squid. This work is so important because, even though they originally used it to explain the propagation of electrical signals in a nerve cell, it has since then been extended and applied to model a wide variety of electrically excitable cells. The model (and equations) they proposed is mentioned, in brief, below.

In the previous section we described how the cell membrane could be modelled as a capacitor, resulting in the following equation:

$$C_m \frac{dV}{dt} + I_{ionic}(V, t) = 0 \quad (1.4)$$

In the giant squid axon, the main ionic currents are the sodium and potassium currents. HH lumped together the other, less significant, ion currents into one current called the *leak current* ( $I_L$ ). HH assumed the membrane channel conductances to have linear I-V curves (this choice was largely dictated by experimental data) and so, we get:

$$C_m \frac{dV}{dt} = -g_{Na}(V - V_{Na}) - g_K(V - V_K) - g_L(V - V_L) \quad (1.5)$$

An important point to note is that the channel conductances are not constant; they seem to depend on the membrane voltage. Determining the conductances by being able to measure individual ionic currents and thus deducing the changes in conductances (with V) was what was brilliantly accomplished by Hodgkin and Huxley. They found, by performing various voltage clamp and space clamp experiments, that:

**For the potassium current:**

$$g_K = \bar{g}_k n^4$$

$$\frac{dn}{dt} = \alpha_n(1 - n) - \beta_n n \quad (1.6)$$

where  $n$  is an *activating* variable.

**For the Sodium current:**

$$g_{Na} = \bar{g}_{Na} m^3 h$$

$$\frac{dm}{dt} = \alpha_m(1 - m) - \beta_m m \quad (1.7)$$

$$\frac{dh}{dt} = \alpha_h(1 - h) - \beta_h h \quad (1.8)$$

where  $m$  is an *activating* variable and  $h$  is an *inactivating* variable.

**The specific forms of  $\alpha$ :**

$$\alpha_m = 0.1 \frac{25 - V}{e^{\left(\frac{25-V}{10}\right)} - 1}$$

$$\alpha_h = 0.07 e^{\left(\frac{-V}{20}\right)}$$

$$\alpha_n = 0.01 \frac{10 - V}{e^{\left(\frac{10-V}{10}\right)} - 1}$$

**The specific forms of  $\beta$ :**

$$\beta_m = 4 e^{\left(\frac{-V}{18}\right)}$$

$$\beta_h = \frac{1}{e^{\left(\frac{30-V}{10}\right)} + 1}$$

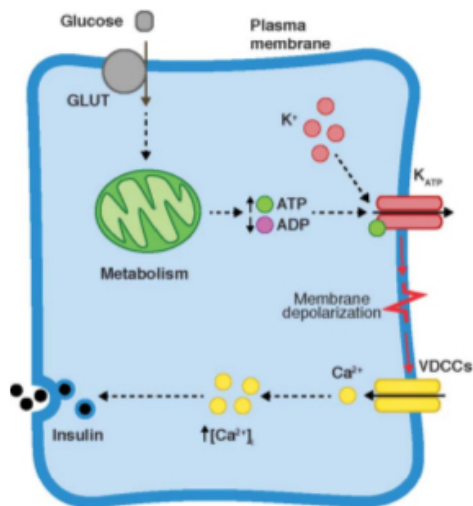
$$\beta_n = 0.125 e^{\left(\frac{-V}{18}\right)}$$

Equations 1.5, 1.6, 1.7 and 1.8 are the four main equations, which, along with the functional forms of the activating/inactivating variables make up the HH model.

All the models describing the electrical activity of beta cells are based on the HH equations.

## 1.3 Glucose Stimulated Insulin Secretion (GSIS) in Beta Cells

Pancreatic beta cells synthesize and secrete the hormone insulin, the only known hormone responsible for lowering blood glucose levels. Under normal conditions, insulin is secreted in a pulsatile manner<sup>1,2</sup>. Interestingly, insulin oscillations are disrupted in patients with Type 2 diabetes<sup>3</sup>. In beta cells, glucose stimulated insulin secretion (GSIS) is driven by a well established sequence of events (*Fig. 1.4*).



Glucose is metabolized after it is transported into the cell by GLUT transporters. This triggers the production of ATP and the subsequent closure of K<sub>ATP</sub> channels. The membrane gets depolarized and VDCCs open, allowing an influx of Ca<sup>+2</sup> ions. The increase in [Ca<sup>+2</sup>]<sub>i</sub> stimulates insulin secretion in a Ca<sup>+2</sup> - dependent manner.

Figure 1.4: *Schematic of GSIS in beta cells*

GSIS begins with the transport of glucose into the cell through glucose transporters (GLUT). The increased intracellular concentration of glucose accelerates metabolism and hence, the production of ATP at the expense of ADP. This results in an increase in the ATP/ADP ratio, causing the inactivation of ATP dependent Potassium channels (K<sub>ATP</sub>), resulting in the slow depolarization of the membrane potential. Once the membrane potential crosses a threshold value, Voltage Dependent Calcium Channels (VDCCs) open and Ca<sup>+2</sup> enters the cell. It is this increase in [Ca<sup>+2</sup>]<sub>i</sub> that cause insulin secretion. This is the main pathway of GSIS (also called the triggering or K<sub>ATP</sub>-dependent pathway).

As a complement to experimental work, mathematical models of beta cells have been built to better understand how various cellular mechanisms involved in GSIS interact and to provide feasible explanations of experimental observations and predictions. As more and more experimental evidence has emerged, models have grown in complexity: from the early, minimal model that just included a few ion channels, to models that incorporate detailed representations of metabolism, glycolysis and electrical activity.

## 1.4 Aim and Organization of the Thesis

The aim of this study is, broadly, to understand how mathematical models are formulated based on experimental information; how models are improved by incorporating new experimental information; how computational techniques can be used to mimic experimental manipulations and study the resultant behaviour of the theoretical model; and finally, to offer new theoretical predictions that can be tested experimentally. Specifically, I have used Pancreatic beta cells as my model.

In chapter 2 all the single cell models and methods used have been described.

In chapter 3 the simulation results of the electrical activity of the single cell models, two coupled cells and multiple coupled cells have been presented.

In chapter 4 a summary of the results and possible future directions the work in this thesis can be extended to have been mentioned.

In chapter 5 the literature referred to in this thesis has been listed chapter-wise.

All the MATLAB codes used have been attached in Appendix A.



# Chapter 2

## Models and Methods

### 2.1 Models

Mathematical models of beta cells have been proposed as tools to describe how various cellular processes interact in GSIS. As models have evolved over time, more and more of the cellular processes (described in the previous section) have been included in models of beta cells.

The pioneering model was published in 1983 by Chay and Keizer<sup>1</sup>. This model builds on the Hodgkin-Huxley paradigm<sup>2</sup>. The first models (Chay-Keizer<sup>1</sup>, Chay<sup>13</sup>) explained fast bursting of membrane potential correctly but predicted oscillations of cytosolic  $[Ca^{+2}]$  with different kinetics than what was seen experimentally. With the discovery of ATP dependent Potassium Channels ( $K_{ATP}$ )<sup>4,5</sup> in 1984, a set of models including  $K_{ATP}$  channels came up (Keizer and Magnus<sup>11,12</sup>, Smolen and Keizer<sup>6</sup>). In order to explain the origins of extremely slow bursting (period of many minutes)<sup>15,16,17</sup>, a wave of models in which the ER was included as a second  $Ca^{+2}$  compartment came up (Bertram *et. al.*<sup>18</sup>, Chay *et.al.*<sup>19</sup>). Some other scientists have proposed an alternative mechanism to explain the wide range in the periodicity of bursting - the idea that the periodicity is based on the interactions between a fast and a slow variable(s) (Bertram *et. al.*<sup>7</sup>, Bertram and Sherman<sup>8</sup>, Bertram *et. al.*<sup>9</sup>). The current generation of models can explain the full range of observed bursting periods.

All the models mentioned so far (and the ones described in the next sections) are based on rodent experimental data. Human pancreatic beta cells differ from rodent beta cells in a few ways (distribution of cells in the islet<sup>20</sup>, the types of ion channels expressed<sup>21,22</sup> and kinetics of insulin exocytosis<sup>23</sup>). Based on these differences (and based on the existing rodent models), mathematical models of human beta cells are currently in the early stages of development.

In this section four of the models based on rodent beta cells are described - mathematically and in terms of the physiological mechanisms involved.

### 2.1.1 Chay-Keizer Model (1983)

Chay and Keizer<sup>1</sup> (CK) developed their pioneering model of the beta cell in 1983. A schematic of the CK model is shown in figure 2.1. The ion channels they considered are: Voltage Dependent Calcium Channels (VDCCs), Voltage Dependent Potassium Channels ( $K_V$ ) and Calcium Dependent Potassium Channels ( $K_{Ca}$ ). The form of the equations used to describe the ion channels is the HH type<sup>2</sup>. Intracellular calcium handling is modelled in a minimal manner. As proposed by Atwater et. al.<sup>3</sup>, the CK model uses the effects of intracellular calcium concentrations ( $[Ca^{+2}]_i$ ) on  $K_{Ca}$  channels as the mechanism to terminate or initiate bursts of potential.

The active phase, shown by (1) in figure 2.1, is sustained by  $K_V$  and VDCCs and  $[Ca^{+2}]_i$  slowly rises. Eventually, the  $K_{Ca}$  channels get activated and the membrane repolarizes (shown by (2)). During the silent phase (shown by (3)), the VDCCs and  $K_V$  channels are closed and  $[Ca^{+2}]_i$  is extruded from the cell, inhibiting the activity of the  $K_{Ca}$  channels. This causes depolarization of the membrane potential, activating the VDCCs and  $K_V$  channels, initiating a new burst.

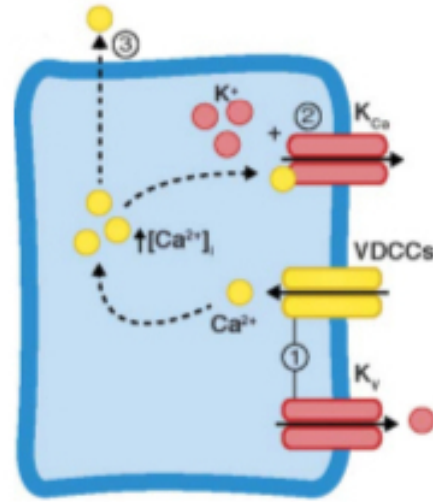


Figure 2.1: Schematic of the Chay-Keizer Model [14]

#### The Model Equations

##### The differential equations:

$$\frac{dV}{dt} = \frac{1}{C_m} (2I_{Ca} + I_K + I_{KCa} + I_l)$$

$$\frac{dm_{Ca}}{dt} = \phi (\alpha_{mCa}(1 - m_{Ca}) + \beta_{mCa}m_{Ca})$$

$$\frac{dh_{Ca}}{dt} = \phi (\alpha_{hCa}(1 - h_{Ca}) + \beta_{hCa}h_{Ca})$$

$$\frac{dn}{dt} = \phi (\alpha_n(1 - n) + \beta_n n)$$

$$\frac{dCa}{dt} = f \left( \frac{3}{r} I_{Ca} - K_{Ca} Ca \right)$$

##### The ionic currents:

$$I_{Ca} = g_{Ca} m_{Ca}^3 h_{Ca} (V_{Ca} - V)$$

$$I_K = g_K n^4 (V_K - V)$$

$$I_{KCa} = g_{KCa} \frac{Ca}{Ca + K_{diss}} (V_K - V)$$

$$I_l = g_l (V_l - V)$$

Parameter	Value
$C_m(\mu\text{F}/\text{cm}^2)$	1
$g_{Ca}(\text{mS}/\text{cm}^2)$	6.5
$g_K(\text{mS}/\text{cm}^2)$	12
$g_l(\text{mS}/\text{cm}^2)$	0.04
$g_{KCa}(\text{mS}/\text{cm}^2)$	0.09
$V_{Ca}(\text{mV})$	100
$V_K(\text{mV})$	-75
$V_l(\text{mV})$	-40
$V'(\text{mV})$	50
$V^*(\text{mV})$	30
$T(^{\circ}\text{C})$	20
$r(\mu\text{m})$	$8.9 \times 10^{-4}$
$f$	0.004
$K_{diss}(\mu\text{M})$	1
$K_{Ca}(\mu\text{M})$	0.04

### Gating functions:

$$\alpha_{mCa} = -0.1 \frac{(V+V'-25)}{e^{-\frac{(V+V'-25)}{10}} - 1}$$

$$\beta_{mCa} = 4e^{-\frac{(V+V')}{18}}$$

$$\alpha_{hCa} = 0.07e^{-\frac{(V+V')}{20}}$$

$$\beta_{hCa} = \frac{1}{e^{-\frac{(V+V')-30}{10}} + 1}$$

$$\alpha_n = -0.01 \frac{(V+V^*-10)}{e^{-\frac{V+V^*-10}{10}} - 1}$$

$$\beta_n = 0.125e^{-\frac{V+V^*}{80}}$$

$$\phi = 3^{\frac{T-6.3}{10}}$$

### 2.1.2 Smolen-Keizer Model (1992)

In 1984 ATP dependent Potassium Channels ( $K_{ATP}$ ) were identified in rodent beta cells <sup>4,5</sup>, emerging as a probably link between metabolism and electrical activity. The activity of  $K_{ATP}$  channels is inhibited by ATP and stimulated by ADP. These channels are responsible for the resting membrane potential of beta cells. Smolen and Keizer <sup>6</sup> (SK) introduced  $K_{ATP}$  channels to the models of beta cells in order to analyze the role of cyclical changes in the ATP/ADP ratio on the electrical activity of beta cells. This model assumes as negative influence of  $\text{Ca}^{+2}$  on the production of ATP. A schematic of the SK model is shown in figure 2.2.

The active phase, shown by (1) in figure 2.2, is sustained by  $K_V$  and VDCCs and  $[\text{Ca}^{+2}]_i$  slowly rises. This exerts a negative effect on the production of ATP (reflected in the increase of ADP) causing a corresponding decrease in the ATP/ADP ratio. As a consequence,  $K_{ATP}$  channels open, repolarizing the membrane (shown by (2)). During the silent phase (shown by (3)), the  $K_V$  channels and VDCCs are closed and  $[\text{Ca}^{+2}]_i$  is extruded from the cell. As  $[\text{Ca}^{+2}]_i$  decreases, the production of ATP is

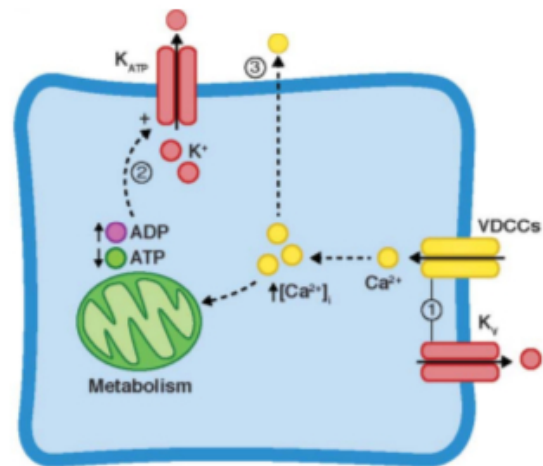


Figure 2.2: Schematic of the Smolen-Keizer Model [14]

potentiated, closing the  $K_{ATP}$  channels and initiating a slow depolarization of the membrane, initiating a new cycle.

### The Model Equations

#### The differential equations:

$$\frac{dV}{dt} = -\frac{1}{C_m} (I_{Ca} + I_K + I_{KATP})$$

$$\frac{dn}{dt} = \lambda_a(\phi_k - n) \frac{1}{\tau_n}$$

$$\frac{dCa}{dt} = f \cdot (\alpha I_{Ca} - a_{kca} Ca - \frac{d_{amp} ca^{hill}}{a_{kdamp}^{hill} + ca^{hill}})$$

$$\frac{dd}{dt} = k_a((tot - d) - d)$$

$$\frac{dpc}{dt} = k_{1p} pot(1 - f_{2\infty}) - k_{1m} pc$$

$$\frac{dpot}{dt} = -(k_{1p} pot \cdot (1 - f_{2\infty}) - k_{1m} pc - \lambda_k(k_{3p} \cdot f_{2\infty} \cdot pot - k_{3m}(1 - pot \cdot pc)))$$

$$\frac{dii}{dt} = \lambda_i \frac{(\phi_i - ii)}{\tau_i}$$

$$\frac{djj}{dt} = \lambda_j \frac{(\phi_j - jj)}{\tau_j}$$

#### The ionic currents:

$$I_{Ca} = g_{Cab} g_{hk}$$

$$I_{KV} = g_k \cdot ii \cdot n(V - V_K)$$

$$I_{KATP} = g_{ATP}(V - V_K)$$

Other equations:

$$\phi_{CaJ} = \frac{1}{1+e^{-\frac{V_{mj}-V}{S_{mj}}}}$$

$$\phi_K = \frac{1}{1+e^{-\frac{V_n-V}{S_n}}}$$

$$\phi_{Ca} = \frac{1}{1+e^{-\frac{V_m-V}{S_m}}}$$

$$\phi_i = \frac{1}{1+e^{-\frac{V_i+V}{S_i}}}$$

$$\phi_j = \frac{1}{1+e^{-\frac{V_j+V}{S_j}}}$$

$$k_{1m} = \lambda_m \phi_{ICa}$$

$$k_{1p} = \lambda_m (1 - \phi_{ICa})$$

$$k_{2p} = -k_{2s} g_{hk}$$

$$g_{Cab} = g_{Cam} (\chi P + (1 - \chi) \phi_{CaJ} j_j) + g_{Cal}$$

$$g_{hk} = \frac{Ca_0 V}{1 - e^{13.35}}$$

$$g_{ATP} = \bar{g}_{ATP} \frac{1 + \frac{d}{k_1}}{1 + \frac{d}{k_1} + \frac{tot-d}{k_2}}$$

$$f_{2\infty} = \frac{k_{2p}}{k_{2m} + k_{2p}}$$

Paramater	Values	Parameter	Values
$V_K$ (mV)	-75	$g_{cam}$ (pS)	1470
$C_m$ ( $\mu\text{F}/\text{cm}^2$ )	5309	$\chi$	0.27
$V_{mj}$ (mV)	-11	$S_{mj}$ (mV)	3.6
$g_{cal}$ (pS)	6.3	$Ca_0$ (mM)	7
$g_K$ (pS)	5000	$\bar{g}_{ATP}$ (pS)	6000
$K_1$ ( $\mu\text{M}$ )	0.45	$K_2$ ( $\mu\text{M}$ )	0.012
$V_n$ (mV)	-20	$\lambda_a$	1.2
$S_n$ (mV)	5.3	$f$	0.03
$k_a$ ( $\text{ms}^{-1}$ )	$2 \times 10^{-5}$	$r$ ( $\mu\text{m}$ )	0.76
$r_1$ ( $\mu\text{m}$ )	0.35	$\lambda_m$	0.78
$V_m$ (mV)	-15	$S_m$ (mV)	6.2
$k_{2s}$ ( $\text{ms}^{-1}$ )	1.512	$k_{2m}$ ( $\text{ms}^{-1}$ )	65
$\lambda_K$	1	$k_{3p}$ ( $\text{ms}^{-1}$ )	0.02
$\lambda_i$	1	$V_i$ (mV)	36
$S_i$ (mV)	4.5	$\tau_i$ (ms)	2600
$\lambda_j$	1	$V_j$	50
$S_j$ (mV)	6.3	$\tau_j$ (ms)	50000

$$\alpha = \frac{3000}{8\pi(96487)^{73}}$$

$$\alpha_n = e^{-\frac{V+75}{65}}$$

$$\beta_n = e^{-\frac{V+75}{65}}$$

$$\tau_n = \frac{60}{\alpha_n + \beta_n}$$

$$\tau_j = \tau_{jmin} + \frac{\tau_j}{e^{\frac{V+V_j}{2S_j}} + e^{-\frac{V+V_j}{2S_j}}}$$

### 2.1.3 Phantom Burster Model (2000)

Bertram *et. al.*<sup>7</sup> developed a model based on the idea that the periodicity of bursting is determined by the interaction between a fast and a slowly oscillating variable. This model is capable of producing bursting with an intermediate period, distinct from that of the fast and the slow variables. Hence, models of this type are called phantom bursters.  $s_1$  is the faster variable and  $s_2$  the slower one. For the purposes of the model it is sufficient if  $I_{s_1}$  and  $I_{s_2}$  are repolarizing, negative feedback currents that turn on

when the cell is depolarized (Alternatively, they could be depolarizing currents that turn off when the cell is depolarized). There are many candidates for the identity of the fast and slow processes - their exact biophysical identities remain obscure. Bertram *et. al.*<sup>7</sup> suggest that  $I_{s1}$  could be thought of as a current through  $K_{Ca}$  channels and  $I_{s2}$  as a current through  $K_{ATP}$  channels. Notice that neither of these channels is directly voltage dependent but both channels respond to changes in  $[Ca^{+2}]_i$  levels which is voltage dependent.

### The Model Equations

#### The differential equations:

$$\frac{dV}{dt} = -\frac{1}{C_m} (I_{Ca} + I_K + I_l + I_{s1} + I_{s2} + I_{stim})$$

$$\frac{dn}{dt} = \lambda \frac{n_{\infty} - n}{\tau_n}$$

$$\frac{ds_1}{dt} = \frac{s_{1\infty} - s_1}{\tau_{s1}}$$

$$\frac{ds_2}{dt} = \frac{s_{2\infty} - s_2}{\tau_{s2}}$$

#### The ionic currents:

$$I_{Ca} = g_{Ca} m_{\infty} (V - V_{Ca})$$

$$I_K = g_K n (V - V_K)$$

$$I_l = g_l (V - V_l)$$

$$I_{s1} = g_{s1} s_1 (V - V_K)$$

$$I_{s2} = g_{s2} s_2 (V - V_K)$$

#### Gating and other functions:

$$m_{\infty} = \frac{1}{1 + e^{\frac{V_m - V}{s_m}}}$$

$$n_{\infty} = \frac{1}{1 + e^{\frac{V_n - V}{s_n}}}$$

$$\tau_n = \frac{\bar{\tau}_n}{1 + e^{\frac{V - V_n}{s_n}}}$$

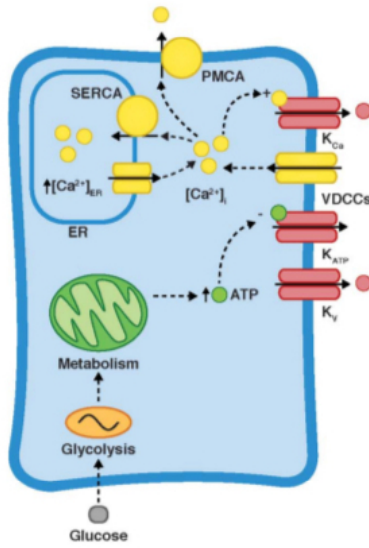
$$s_{1\infty} = \frac{1}{1 + e^{\frac{V_{s1} - V}{s_{s1}}}}$$

$$s_{2\infty} = \frac{1}{1 + e^{\frac{V_{s2} - V}{s_{s2}}}}$$

Parameter	Values	Parameter	Values
$V_l$ (mV)	-40	$g_{s1}$ (pS)	32
$V_{Ca}$ (mV)	100	$g_{s2}$ (pS)	3
$V_K$ (mV)	-80	$C_m$ ( $\mu\text{F}/\text{cm}^2$ )	4524
$\lambda$	1.1	$\bar{\tau}_n$ (ms)	9.09
$g_{Ca}$ (pS)	280	$V_m$ (mV)	-22
$g_K$ (pS)	1300	$V_n$ (mV)	-9
$g_l$ (pS)	25	$s_m$	7.5
$V_{s1}$ (mV)	-40	$s_n$	10
$\tau_{s1}$ (ms)	1000	$s_{s1}$	0.5
$V_{s2}$ (mV)	-42	$s_{s2}$	0.4
$\tau_{s2}$ (ms)	$1.2 \times 10^5$		

### 2.1.4 Dual Oscillator Model (2007)

In 2004, Bertram and Sherman<sup>8</sup> proposed a model using the phantom bursting mechanism with 3 slow processes ( $[Ca^{+2}]_i$ ,  $[Ca^{+2}]_{ER}$ , and ATP/ADP). In a later model, called the Dual Oscillator Model<sup>9</sup> (DOM), they combined this model with a model of glycolysis<sup>10</sup> and a model for mitochondrial metabolism<sup>11,12</sup>. A schematic of the DOM is shown in figure 2.3.



The interactions between the glycolytic, metabolic and electrical components drive different electrical behaviours depending on the regime of the glycolytic and electrical components. ATP production is controlled by the rate at which glucose is metabolized. Changes in ATP levels mediates the conductance of  $K_{ATP}$  channels, depolarization and  $Ca^{+2}$  influx. All three compartments are affected by changes in  $[Ca^{+2}]_i$ .

Figure 2.3: *Schematic of the Dual Oscillator Model [14]*

### The Model Equations

#### The differential equations:

$$C_m \frac{dV}{dt} = -(I_K + I_{Ca} + I_{K(Ca)} + I_{K(ATP)})$$

$$\frac{dn}{dt} = \frac{n_\infty(V) - n}{\tau_n}$$

$$\frac{dCa}{dt} = f_{cyt}(J_{mem} + J_{er})$$

$$\frac{dCa_{er}}{dt} = -f_{er}\left(\frac{V_{cyt}}{V_{er}}\right)J_{er}$$

$$\frac{dG6P}{dt} = \lambda(R_{GK} - R_{PFK})$$

$$\frac{dFBP}{dt} = \lambda\left(R_{PFK} - \frac{1}{2}R_{GPDH}\right)$$

$$\frac{dATP}{dt} = \frac{ATP - ADP e^{[(r+\gamma)(1-ca/r_1)]}}{\tau_a}$$

#### The ionic currents:

$$I_K = \bar{g}_K n (V - V_K)$$

$$I_{Ca} = \bar{g}_{Ca} m_\infty(V) (V - V_{Ca})$$

$$I_{K(Ca)} = g_{K(Ca)} (V - V_K)$$

$$I_{K(ATP)} = g_{K(ATP)} (V - V_K)$$

**Cytosolic and ER calcium dynamics:**

$$J_{mem} = \alpha I_{Ca} + k_{PMCA} Ca$$

$$J_{er} = J_{leak} - J_{SERCA}$$

$$J_{leak} = p_{leak}(Ca_{er} - Ca)$$

$$J_{SERCA} = K_{SERCA} Ca$$

**Other functions:**

$$g_K(Ca) = \bar{g}_K(Ca) \frac{Ca^2}{K_D^2 + Ca^2}$$

$$g_K(ATP) = \bar{g}_K(ATP) o_\infty(ADP, ATP)$$

$$n_\infty(V) = \frac{1}{1 + e^{-(16+V)/5}}$$

$$m_\infty(V) = \frac{1}{1 + e^{-(20+V)/12}}$$

$$o_\infty(ADP, ATP) = \frac{0.08 \left(1 + \frac{2MgADP^-}{17\mu M}\right) + 0.89 \left(\frac{MgADP^-}{17\mu M}\right)}{\left(1 + \frac{MgADP^-}{17\mu M}\right)^2 \left(1 + \frac{ADP^{3-}}{26\mu M} + \frac{ATP^{4-}}{1\mu M}\right)}$$

**Nucleotide concentrations:**

$$A_{tot} = AMP + ADP + ATP$$

$$AMP = \frac{ADP^2}{ATP}$$

$$MgADP^- = 0.165ADP$$

$$ADP^{3-} = 0.135ADP$$

$$ATP^{4-} = 0.05ATP$$

**Glycolysis:**

$$w_{ijkl} = \frac{1}{f_{13}^{ik} f_{23}^{jk} f_{41}^{il} f_{42}^{jl} f_{43}^{kl}} \left(\frac{AMP}{K_1}\right)^i \left(\frac{FBP}{K_2}\right)^j \left(\frac{F6P}{K_3}\right)^k \left(\frac{ATP^2}{K_4}\right)^l$$

$$R_{PFK} = V_{max} \frac{(1-\lambda)w_{1110} + \lambda \sum_{ijl} w_{ijl1}}{\sum_{ijkl} w_{ijkl}}$$

$$F6P = 0.3G6P$$

$$R_{GPDH} = 0.2\sqrt{FBP}$$

Parameter	Value	Parameter	Value	Parameter	Value
$\lambda$	0.06	$K_1$ ( $\mu M$ )	30	$K_2$ ( $\mu M$ )	1
$K_3$ ( $\mu M$ )	$5 \times 10^4$	$K_4$ ( $\mu M$ )	220	$f_{23}$	0.2
$f_{41}$	20	$f_{42}$	20	$f_{43}$	20
$C$ (fF)	5300	$\tau_n$ (ms)	20	$\bar{g}_K$ (pS)	2700
$\bar{g}_{Ca}$ (pS)	1000	$\bar{g}_{KATP}$ (pS)	$2.5 \times 10^4$	$V_K$ (mV)	-75
$V_{Ca}$ (mV)	25	$K_D$ ( $\mu M$ )	0.5	$\alpha$ ( $\mu M/ms$ )	$4.5 \times 10^{-6}$
$k_{pmca}$ ( $ms^{-1}$ )	0.1	$p_{leak}$ ( $ms^{-1}$ )	$2 \times 10^{-4}$	$k_{SERCA}$ ( $ms^{-1}$ )	0.4
$f_{er}$	0.01	$f_c$	0.01	$\frac{V_{Cyt}}{V_{er}}$	31
$V_{max}$ (mV)	2	$A_{tot}$ ( $\mu M$ )	3000	$v_\gamma$	2.2
$k_\gamma$ ( $\mu M/s$ )	10	$R_{GK}$ ( $s^{-1}$ )	0.2	$K_{dd}$ ( $\mu M$ )	17
$K_{td}$ ( $\mu M$ )	26	$K_{tt}$ ( $\mu M$ )	1	$\tau_a$ (ms)	$3 \times 10^5$
$r$	1	$r_1$ ( $\mu M$ )	0.35	$\bar{g}_K(Ca)$ (pS)	600



## 2.2 Methods

### 2.2.1 Coupled Map Lattices to Model Islets

Coupled Map Lattices is a method to model the collective dynamics of a system with interacting parts. I have modelled the pancreatic islet as a 1D ring of cells (with periodic boundary conditions, *Fig. 2.4*). Each cell is coupled to its nearest neighbour. Biologically, coupling is achieved through the presence of gap junctions in the membranes of adjoining cells which allows small molecules and ions to pass through. Hence, neighbouring cells can sense voltage differences through the gap junctions. In the 1D ring, each lattice site (or cell) is described by the DOM system of equations.

The effect of coupling is modelled by changing the form of the voltage differential equation. Equation (1.4) is modified such that the voltage variation of the  $i^{\text{th}}$  cell can be written as:

$$C_m \frac{dV_i}{dt} = -\sum I_{ionic} - g_c(V_{i-1} - V_i) - g_c(V_i - V_{i+1})$$

where,  $g_c$  is the gap junction coupling strength between two adjacent cells and  $V_{i+1}$  and  $V_{i-1}$  are the voltages of the two nearest neighbours of the  $i^{\text{th}}$  beta cell.

### 2.2.2 Synchronization Order Parameter, R

The Synchronization Order Parameter <sup>13</sup>,  $R$ , is a measure of spatio-temporal synchronicity in a population of coupled oscillators (here beta cells). The value of  $R$  indicates the level of synchronization by comparing the average of the local signal (here membrane voltage of a beta cell) to the global behaviour.  $R$  is explicitly defined as follows:

$$R = \frac{\langle M^2 \rangle - \langle M \rangle^2}{\left[ \langle z_i^2 \rangle - \langle z_i \rangle^2 \right]}$$

where  $[\ ]$  denotes spatial averaging and  $\langle \ \rangle$  denotes temporal averaging.  $M$  is the spatial average of  $z_i$  (voltage values in the case of coupled beta cells) over  $N$  cells at every time point.  $R$  takes up values between 0 and 1. If all the beta cells are completely synchronized with respect to their electrical activity, then the average of the local signal will be the same as the global behaviour. Hence,  $R \sim 1$ . If the beta cells are unsynchronized, the individual states will be different and hence  $R \sim 0$ .

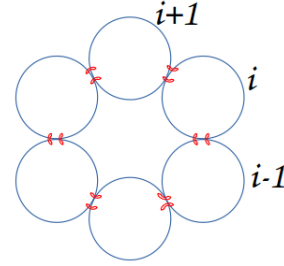


Figure 2.4: 1D ring of coupled cells

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All simulations were programmed on MATLAB R2012b using the ode15s ode solver, and run on a Workstation with an Intel® Xeon(R) CPU E5-1620 0 @ 3.60GHz  $\times$  8 processor. The operating system used was Ubuntu v14.04 LTS (64-bit). It took approximately 5 days to run a  $10^7$ ms simulation of coupled beta cells.

# Chapter 3

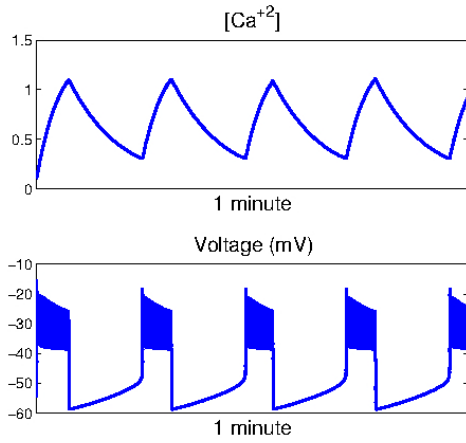
## Results

The results are organized in the following manner: First, the simulation results of the single beta cell bursting dynamics for the four models mentioned in the Models section of Chapter 2 are shown, and their drawbacks mentioned. Second, since the Dual Oscillator Model (DOM) is the only model capable of reproducing most experimental results, simulation results of two cells (each described by the DOM) coupled through gap junctions are shown. Finally, simulation results for a ring-like, one-dimensional CML model of an islet with multiple beta cells, each described by the DOM, coupled to their nearest two neighbours are shown. For both the two and multiple cell models, simulations of a 1D ring of both homogeneous and heterogeneous beta cells (in terms of bursting dynamics) is shown.

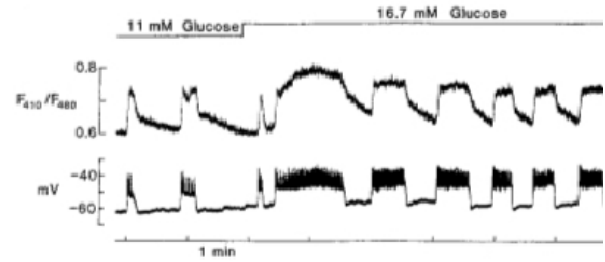
### 3.1 Single Cell Results

#### 3.1.1 Chay-Keizer Model Predictions

The burst pattern and calcium oscillations simulated using the CK model equations (see section 2.1.1) are shown in figure 3.1(a), for the following initial conditions:  $v_0 = -54.774$ ;  $n_0 = 0.00044035$ ;  $m_{ca0} = 0.027532$ ;  $h_{ca0} = 0.086321$ ;  $ca_0 = 0.10749$  (All the values for each of the parameters used are mentioned in the previous chapter). The experimentally observed time curves are shown in figure 3.1(b). It is clear that the CK model predicts a slower and different form of the calcium dynamics than observed experimentally.



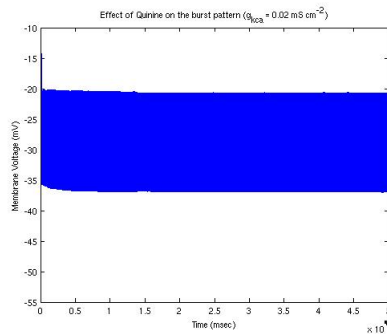
(a) Top: CK predicts the  $[Ca^{+2}]_i$  time curve to be square wave. Bottom: CK predicts the voltage time curve to be square wave with bursting. ( $KCa = 0.04$ ;  $g_{KCa} = 0.09$ )



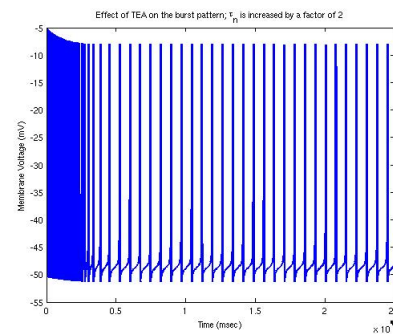
(b) Top: Experimentally observed  $[Ca^{+2}]_i$  time curve. Bottom: Experimentally observed voltage time curve [1]

Figure 3.1: (a) CK model predicted time curves (b) Experimentally observed time curves.

Moreover, it was experimentally found that blocking  $K_{Ca}$  channels with charybdotoxin produced no significant effect on the electrical activity<sup>3</sup>. However, modelling the blocking of  $K_{Ca}$  channels leads to a disruption in the burst pattern (Fig. 1.2 (a)).



(a) CK model predicts a loss in burst pattern on blockage of  $K_{Ca}$  channels by charybdotoxin/quinine. ( $KCa = 0.001$ ;  $g_{KCa} = 0.09$ )



(b) CK model predicts an increase in the voltage dynamics when  $K_V$  channels are blocked by TEA. ( $KCa = 0.04$ ;  $g_{KCa} = 0.02$ )

Figure 3.2: (a) CK model predicted burst pattern when  $K_{Ca}$  channels are blocked (b) CK model predicted burst pattern when  $K_V$  channels are blocked by TEA

The hypothesis proposed by the CK model was discarded when  $[Ca^{+2}]_i$  was measured (experimentally)<sup>1,2</sup> in beta cells, revealing different and more rapid dynamics than predicted by the CK model.

### 3.1.2 Smolen-Keizer Model Predictions

The burst pattern and calcium oscillations simulated using the SK model equations (see section 2.1.2) are shown in figure 3.3, for the following initial conditions:  $v_0 = -53$ ;  $n_0 = 0.002$ ;  $ca_0 = 0.15$ ;  $d_0 = 0.42$ ;  $pc_0 = 0.9$ ;  $pot_0 = 0.02$ ;  $ii_0 = 0.98$ ;  $jj_0 = 0.47$  (All the values for each of the parameters used are mentioned in the previous chapter).

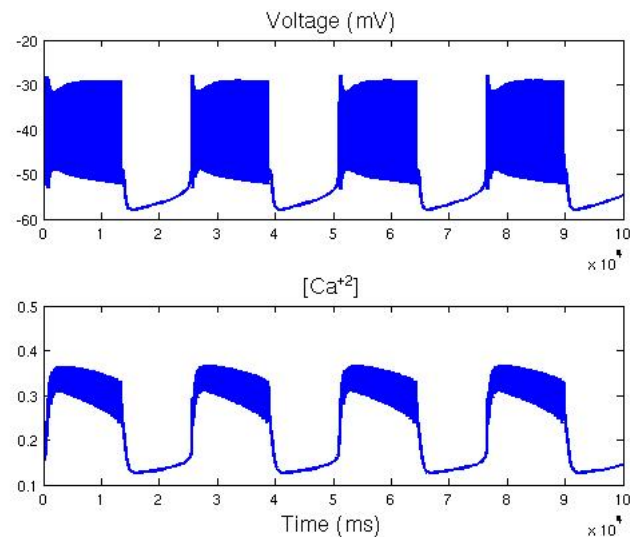
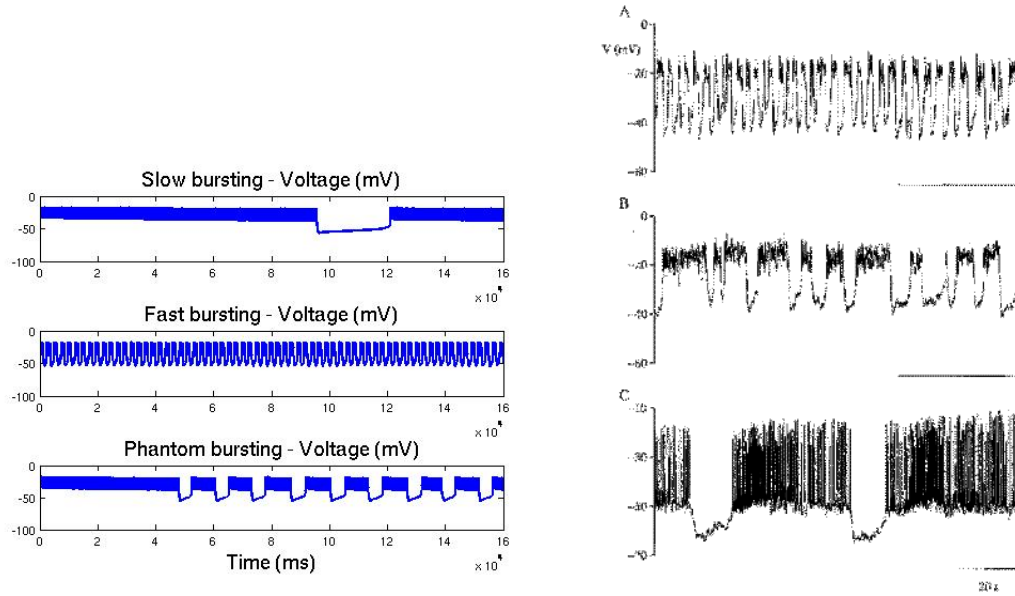


Figure 3.3: *The SK model is able to reproduce the fast dynamics of  $[Ca^{+2}]_i$  oscillations*

Dufer *et. al.*<sup>4</sup> report that electrical activity in beta cells lacking functional  $K_{ATP}$  channels persists. All the models (including the SK model) that are based on the oscillations of the ATP/ADP ratio to produce bursting electrical activity by regulating the conductance of the  $K_{ATP}$  channels are unable to reproduce this observation. However, the SK model is able to reproduce the fast dynamics of  $[Ca^{+2}]_i$  oscillations, and in fact, includes an improved model of the  $Ca^{+2}$  currents (*Fig. 1.2 (b)*).

### 3.1.3 Phantom Burster Model Predictions

The possible burst patterns simulated using the PB model equations (see section 2.1.3) are shown in figure 3.4(a), for the following initial conditions:  $v_0 = -43.0$ ;  $n_0 = 0.03$ ;  $s_{10} = 0.1$ ;  $s_{20} = 0.434$  (All the values for each of the parameters used are mentioned in the previous chapter). The experimentally observed burst patterns are shown in figure 3.4(b). In the model the value for 'gs1' is changed to see bursting with different time periods.  $gs1 = 3$  for slow bursting;  $gs1 = 4$  for intermediate, 'phantom bursting';  $gs1 = 20$  for fast bursting.



(a) The PB model can reproduce slow (top,  $gsI=3$ ), fast (middle,  $gsI=20$ ) and phantom (bottom,  $gsI=4$ ) voltage dynamics.

(b) Experimentally observed burst patterns on fast (A), medium (B), and slow (C) time scales [9]

Figure 3.4: (a) PB model predicted possible voltage dynamics (b) Experimentally observed voltage dynamics

Single cell electrophysiological studies<sup>5,6</sup> have established that single cells can burst as well as spike repetitively, but this bursting is typically much faster than that in islets. Sometimes bursts much slower than those in islets (period 1–6 min)<sup>7,8</sup> are seen. Bursting with a period comparable to that in islets is seen, but only rarely. The Phantom Burster model is capable of reproducing bursting of all these periodicities. The phantom bursting is significant in light of the fact that no slow process with a medium-scale time constant has been identified in beta cells.

### 3.1.4 Dual Oscillator Model Predictions

The burst pattern and calcium oscillations simulated using the DOM equations (see section 2.1.4) are shown in figure 3.5(a), for the following initial conditions:  $V_0 = -60$ ;  $n_0 = 0$ ;  $Ca_0 = 0.1$ ;  $CaER_0 = 185$ ;  $ADP_0 = 780$ ;  $G_6P_0 = 200$ ;  $FBP_0 = 40$  (All the values for each of the parameters used are mentioned in the previous chapter).

The different types of bursting are reproduced by changing the values of  $R_{GK}$  (Rate of the enzyme *Glucokinase*),  $g_{KATP}$  (conductance of the ATP dependent potassium channel) and  $g_{KCa}$  (conductance of the calcium dependent potassium channel) values.

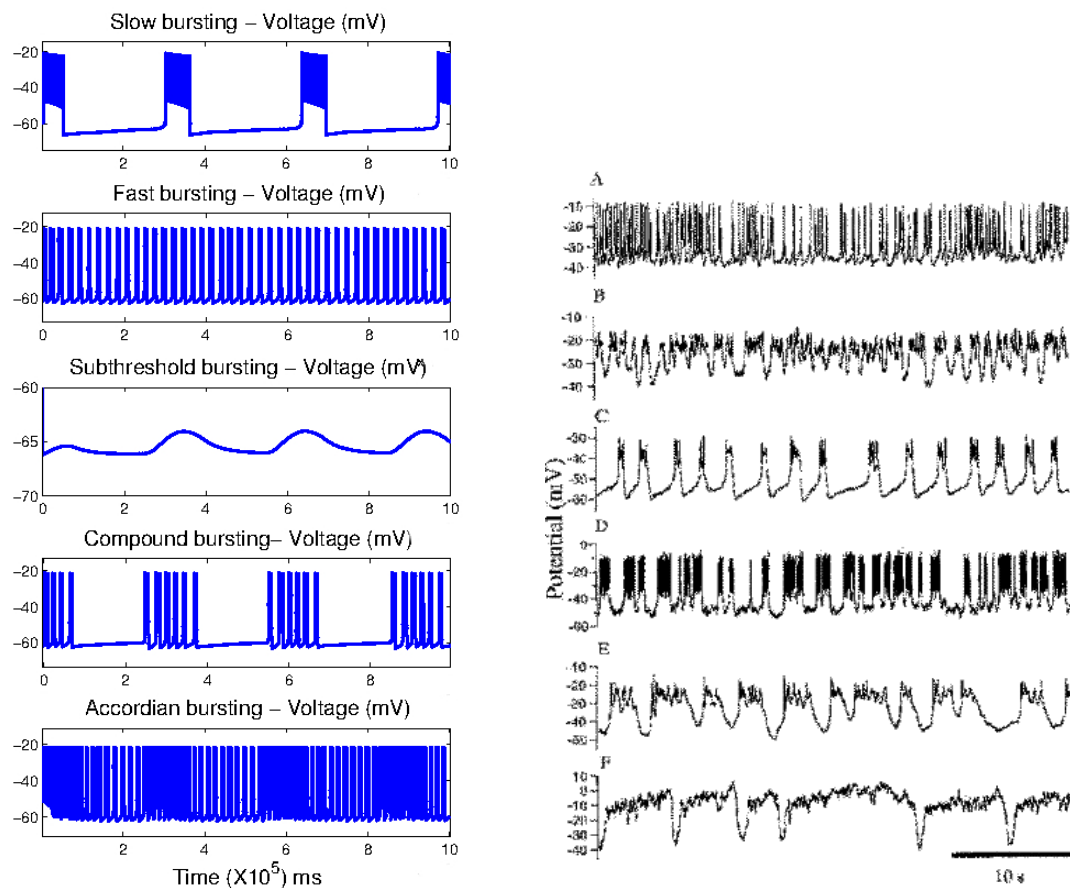
For Slow bursting:  $R_{GK} = 0.2$ ;  $g_{KATP} = 27000$ ;  $g_{KCa} = 100$ ;

For Fast bursting:  $R_{GK} = 0.4$ ;  $g_{KATP} = 25000$ ;  $g_{KCa} = 600$ ;

For Subthreshold bursting:  $R_{GK} = 0.2$ ;  $g_{KATP} = 30000$ ;  $g_{KCa} = 100$ ;

For Compound bursting:  $R_{GK} = 0.2$ ;  $g_{KATPbar} = 25000$ ;  $g_{KCabar} = 600$ ;

For Accordion bursting:  $R_{GK} = 0.2$ ;  $g_{KATP} = 23000$ ;  $g_{KCa} = 600$ ;



(a) DOM reproduces the full repertoire of bursting 'types' for different  $R_{GK}$ ,  $g_{KATP}$  and  $g_{KCa}$  values.

(b) Experimental recordings of different burst patterns in single mouse beta cells [5]

Figure 3.5: (a) DOM predicted repertoire of bursting (b) Experimental observations

The DOM reproduces the full range of periods observed in electrical bursting activity (including the compound and mixed oscillations), seen experimentally<sup>5</sup>. Most of the predications of the DOM have acquired experimental verification. For example; the DOM predicts oscillations in the glycolytic pathway, which has been seen in direct experimentation<sup>10</sup>. The DOM is the only model capable of reproducing the following experimental observation: sub-threshold metabolic oscillations in the absence of  $[Ca^{+2}]_i$  oscillations<sup>11</sup> (Fig. 3.4 (a)).

A point to note is that, like every other model, the DOM cannot explain the observation made by Dufer *et. al.*<sup>4</sup>, that the electrical activity in beta cells lacking functional  $K_{ATP}$  channels persists.

The DOM is a really elegant model because of the incorporation of glycolytic, metabolic and electrical components of the beta cell. There is a clear connection between the electrical activity of the membrane and the global function of the beta cell in the secretion of insulin. Since almost all the predications of the DOM have experimental support, the Dual Oscillator Model was used to simulate a collection of coupled beta cells - the *Islet of Langerhans*.

## 3.2 Coupled Cell Results

Simulation results of beta cells (described by the DOM) coupled via gap junctions are shown. All the constituent cells in the coupled system can have the same dynamics (Homogeneous), or can be a mix of more than one type of bursting dynamics (Heterogeneous).

### 3.2.1 Homogeneous Cells

First, how coupling two cells (via gap junctions) may change the individual bursting behaviour of each cell is shown. Second, the emergent, collective behaviour of an islet is shown by coupling 30 beta cells in a 1D ring with periodic boundary conditions.

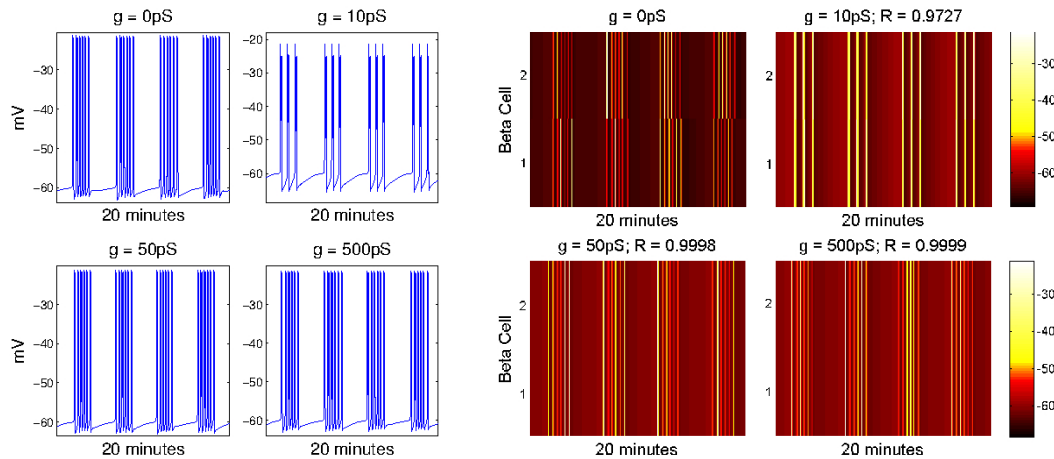
#### 3.2.1.1 Two Cells

The voltage differential equation for two coupled cells is:

$$\text{For Cell 1: } C_m \frac{dV_1}{dt} = -I_{ionic} - g_c(V_1 - V_2)$$

$$\text{For Cell 2: } C_m \frac{dV_2}{dt} = -I_{ionic} - g_c(V_2 - V_1)$$

The burst frequency and length of burst stabilizes for coupling strengths  $> 100\text{pS}$  when two identical cells are coupled.

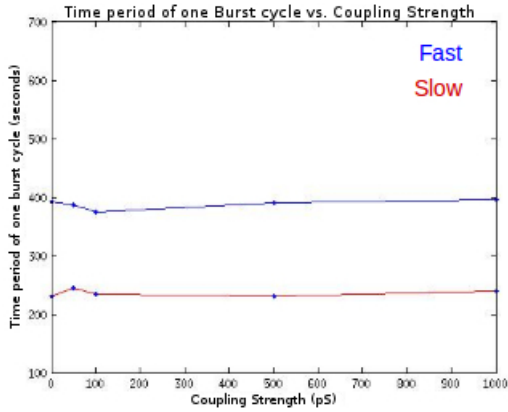


(a) Voltage time curves (for different  $g_c$ s)

(b) Corresponding space-time plots with the synchronisation order parameter,  $R$ , given for coupled cells.

Figure 3.6: (a) Voltage time curves (for different  $g_c$ s) for two coupled compound cells (b) Corresponding space-time plots





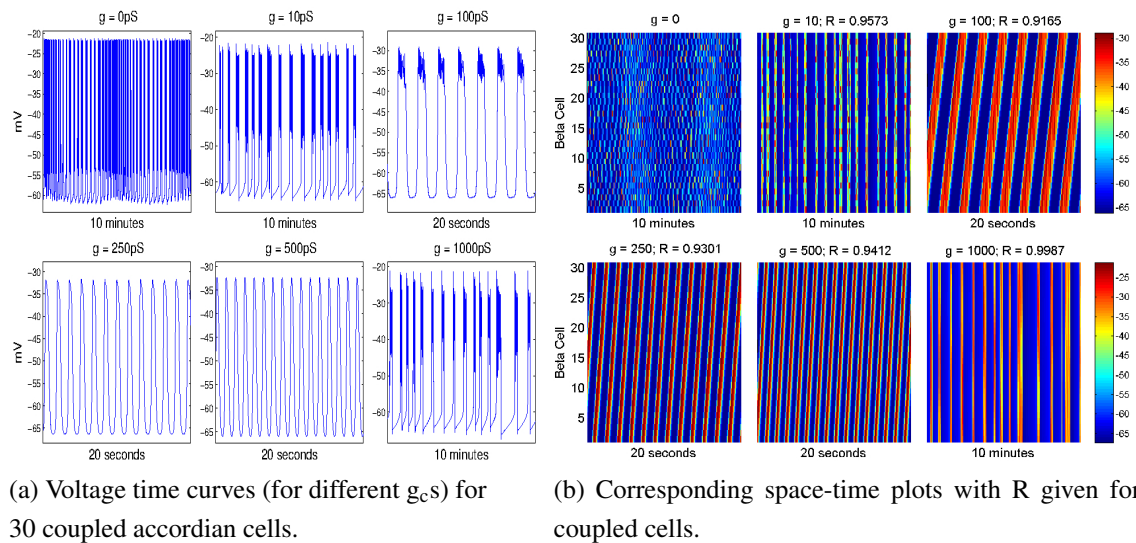
(a) Time period of one cycle (Burst followed by silent phase) versus coupling strength,  $g_c$ .

As can be seen in figure 3.7(a), the time period of one cycle (burst followed by the silent phase) stabilizes when the coupling strength increases beyond  $\sim 100\text{pS}$ . In addition, the percentage of the time period spent bursting also stabilizes. This is true when any two identical (fast, slow, accordian or compound) cells are coupled together.

Figure 3.7: *Time period of one cycle stabilizes for  $g_c > 100\text{pS}$*

### 3.2.1.2 Thirty Cells

On coupling thirty homogeneous cells the burst pattern, time period and burst percentage is lost when  $g_c$  is increased. For intermediate  $g_c$  values ( $\sim 100\text{pS}$  to  $500\text{pS}$ ) the burst pattern is completely lost and replaced by high frequency oscillations.



(a) Voltage time curves (for different  $g_c$ s) for 30 coupled accordian cells.

(b) Corresponding space-time plots with R given for coupled cells.

Figure 3.8: *Voltage time curves and space-time plots for 30 coupled accordian cells*

The specific pattern of the space-time plots for intermediate  $g_c$  values ( $\sim 100\text{pS}$ ,  $250\text{pS}$  and  $500\text{pS}$ ) indicates that the high frequency oscillations of individual beta cells is actually a wave travelling through the ring of beta cell. The wave travels through the lattice with uniform velocity.

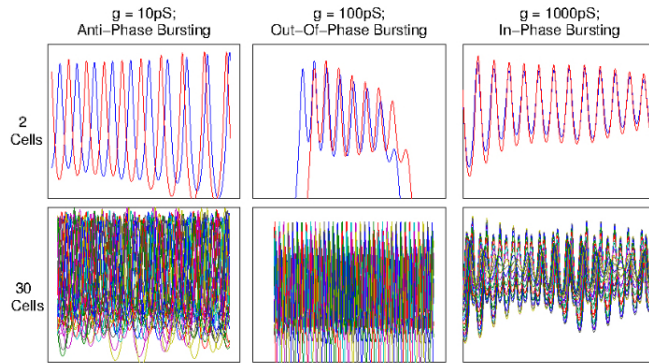
The table in figure 3.9 shows how the time period of one cycle and the percentage of the time period spent on bursting changes as  $g_c$  is increased (for each of the different

cell types). For each homogeneous 1D ring (and each  $g_c$  value) multiple simulations of length between  $10^7$  and  $10^8$  ms were run and the time period values were calculated after a suitable period of transient time was allowed to pass.

	$g = 0\text{pS}$	$g = 10\text{pS}$	$g = 10^2\text{pS}$	$g = 250\text{pS}$	$g = 500\text{pS}$	$g = 10^3\text{pS}$
<b>Accordian</b>	~4.6 mins	~4mins	2 secs (~75%)	2 secs (Osci.)	1.5 secs (Osci.)	Variable (~25%)
<b>Compound</b>	5mins (~60%)	5 mins (~60%)	2 secs (~90%)	2 secs (Osci.)	2 secs (Osci.)	5.2 mins (~65%)
<b>Fast</b>	23 secs (~67%)	54 secs (~80%)	3 secs (~70%)	1.5 secs (Osci.)	2 secs. (Osci.)	Variable (~18%)
<b>Slow</b>	40 secs (~70%)	Variable (~19%)	3 secs (~50%)	2 secs (Osci.)	2 secs (Osci.)	57 secs (~83%)

Figure 3.9: Time period of one cycle and % of time period spent on bursting changes for increasing  $g_c$ .

The cells synchronize their burst and silent phases for low and high coupling strengths. However, for low coupling strengths the bursting time course of neighbouring cells is  $180^\circ$  out of phase. As coupling strength increases, bursting becomes more in phase. For very high coupling strengths bursting is completely in phase.



In figure 3.10, the top panel shows the superimposed time curves of two neighbouring cells (for different  $g_c$ s). The bottom panel shows the superimposed time curves of all 30 cells in the 1D ring. As can be seen, the variation between the time curves is less for a very high  $g_c$  value.

Figure 3.10: Synchronization of bursts goes from  $180^\circ$  anti-phase (low  $g_c$ ), to out-of-phase, to completely in-phase (high  $g_c$ ).

### 3.2.2 Heterogeneous Cells

First, how coupling two different cells may lead to emergent bursting behaviour of a different type is shown. Second, the emergent, collective behaviour of a heterogeneous islet is shown by coupling 32 beta cells in a 1D ring with periodic boundary conditions.

3.2.2.1 Two Cells

Emergent accordian type behaviour is seen on coupling a slow and a compound cell or a fast and a compound cell (Fig. 3.11).

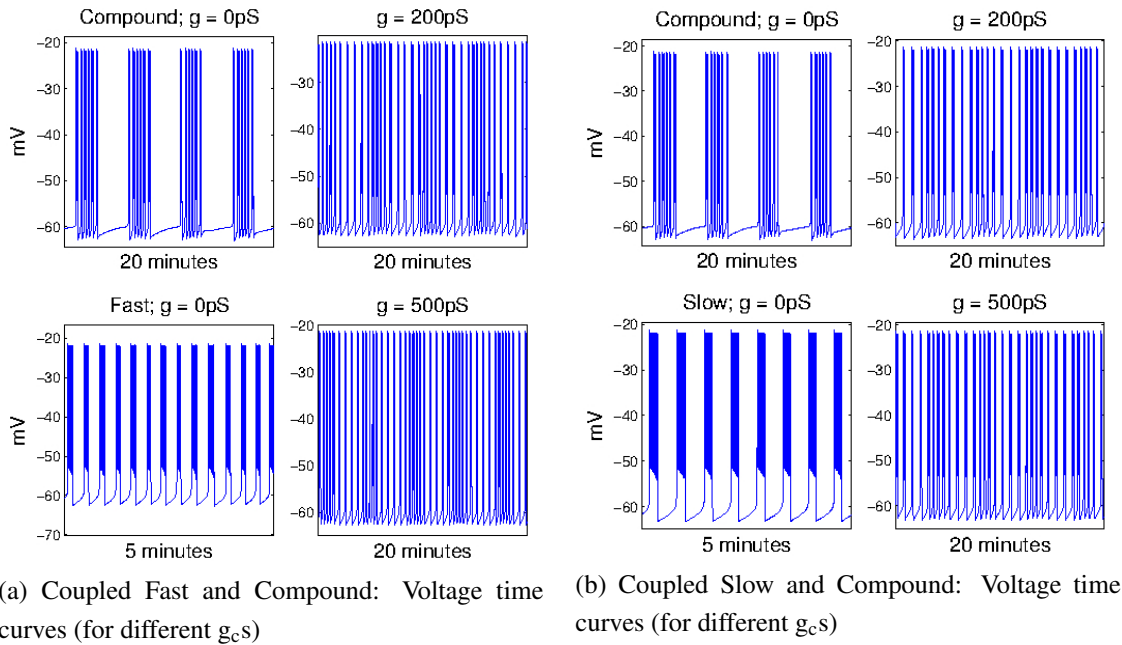


Figure 3.11: Emergent accordian behaviour is seen on coupling a fast/slow cell with a compound cell.

Coupling any other cell types leads to behaviour intermediate to that of the original two cell types.

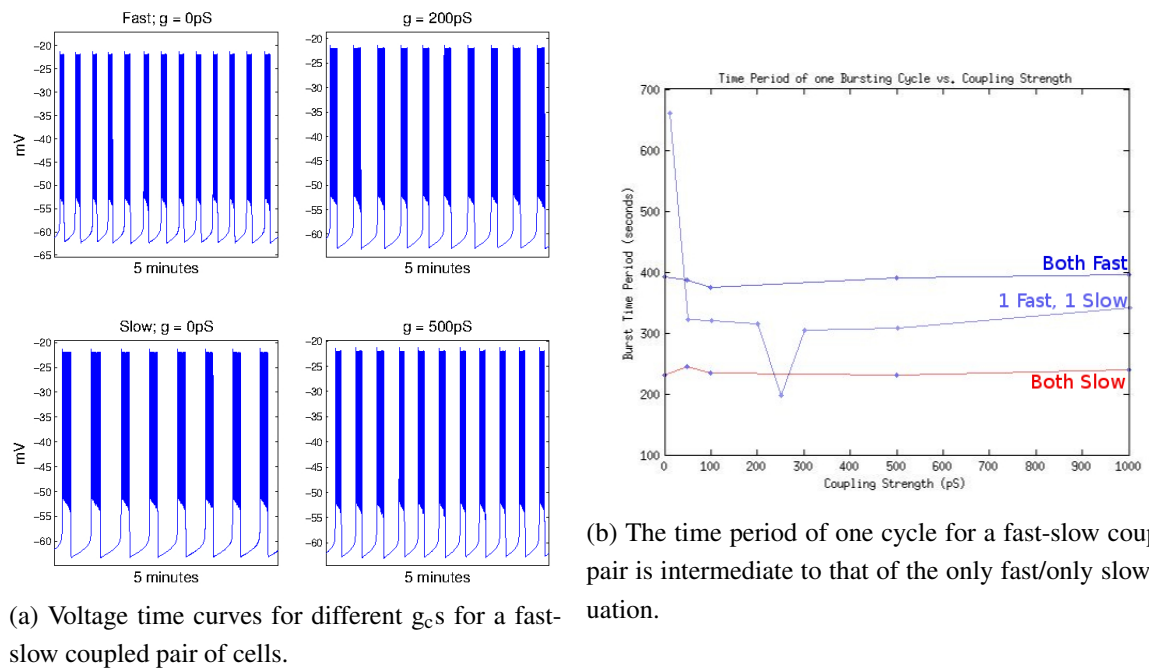
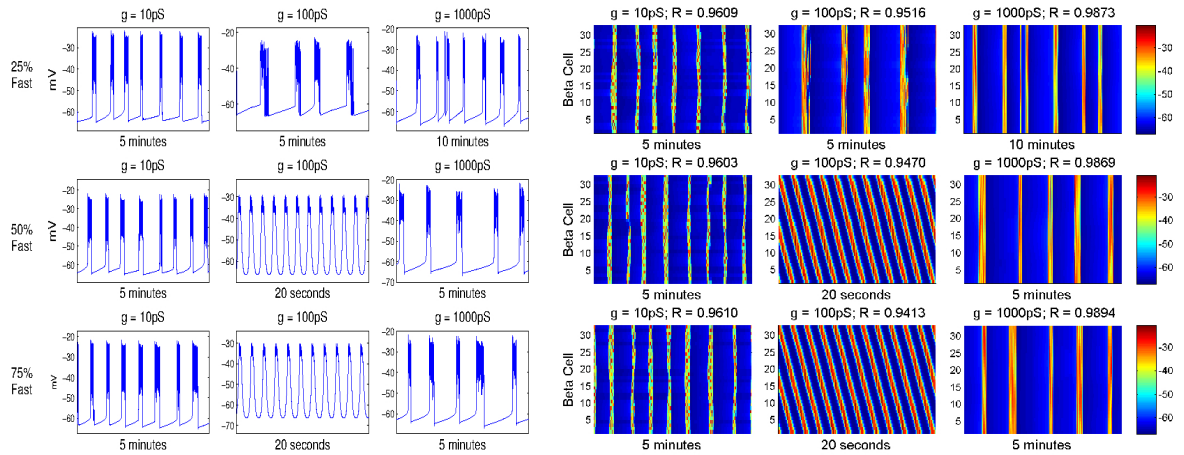


Figure 3.12: Intermediate behaviour of a fast-slow coupled pair of cells.

### 3.2.2.2 Thirty-Two Cells

For large number of cells, coupling of compound and fast/slow types does not lead to emergent accordian type bursting behaviour – instead, the behaviour of such a ring of cells is similar to that of a homogeneous ring of accordian cells (*Fig. 3.13*).



(a) Voltage time curves for 32 coupled cells with (b) Corresponding space-time plots with  $R$  given for covarying percentages of compound and fast cells.      pled cells.

Figure 3.13: *Coupling large number of compound and fast/slow cells doesn't lead to emergent accordian behaviour.*

As the percentage of fast cells increases to 100%, the burst time percentage approaches that of the burst percentage of a homogeneous ring of Fast cells.

The table in figure 3.14 shows how the percentage of the time period spent on bursting changes as  $g_c$  is increased and the percentage of fast cells in a mixture of fast and compound cells is changed. For each heterogeneous 1D ring (and each  $g_c$  value) multiple simulations of length between  $5 \times 10^6$  and  $10^8$  ms were run and the values were calculated after a suitable period of transient time was allowed to pass.

	<b><math>g=0\text{pS}</math></b>	<b><math>g=10\text{pS}</math></b>	<b><math>g=10^2\text{pS}</math></b>	<b><math>g=10^3\text{pS}</math></b>
<b>0% Fast</b>	~40%	~40%	~50%	~36%
<b>25% Fast</b>	-	~17%	~30%	~19%
<b>50% Fast</b>	-	~19%	~25%	~15.4%
<b>75% Fast</b>	-	~20%	~27.5%	~14.8%
<b>100% Fast</b>	~33.3%	~20%	~33.3%	~15%

Figure 3.14: *Burst percentage changes with changing coupling strength and percentage of fast cells in a mixture of fast and compound cells.*

# Chapter 4

## Concluding Remarks

When single units interact, often what is seen is that the behaviour of such a collection of interacting units is different from the behaviour of each individual unit - emergent behaviour. Such is the case with pancreatic beta cells. The dynamics of a single beta cell is extremely different from that of a cluster of coupled cells. The presence of gap junctions between neighbouring beta cells in an islet forces each individual cell to 'talk' to its neighbouring cells. As a consequence, all the cells in the islet behave in a synchronized manner. This synchronization is absolutely essential in the proper functioning of the pancreas; to properly respond to elevated blood glucose levels by secreting insulin in a pulsatile fashion.

Synchronization of the electrical (hence insulin secreting) activity on coupling beta cells means that the body is more tolerant to the presence of a greater variety of beta cells in different 'states' (initial conditions and parameter values). Through this study I have tried to understand how emergent, synchronized behaviour is achieved on coupling cells.

## Conclusions

1. When two homogeneous cells are coupled together (both accordian or compound or fast or slow), the time period of one cycle (burst followed by silent phase) and percentage of time spent bursting stabilizes for  $g_c$  values greater than  $\sim 100\text{pS}$ . No emergent behaviour is seen.
2. When two cells, one compound and the other fast or slow, are coupled together emergent accordian type behaviour is seen. This emergent accordian pattern has the same time period ( $\sim 300$  seconds) as the 'original' accordian type cell. However, the emergent accordian pattern is less dense than the original one.
3. When multiple homogeneous cells are coupled together the original burst pattern of a single cell is lost. For low coupling strengths ( $g_c < \sim 200\text{pS}$ ) the burst

pattern becomes irregular. For intermediate coupling strengths ( $\sim 250\text{pS} < g_c < \sim 500\text{pS}$ ) bursting behaviour is lost and replaced by high frequency (1.5-2 seconds) oscillations. For high coupling strengths ( $g_c > 1000\text{pS}$ ) the irregular bursting pattern is recovered.

4. On observing the space-time plots ( $\sim 250\text{pS} < g_c < \sim 500\text{pS}$ ) of multiple homogeneous coupled cells, it appears that the high frequency oscillations are actually travelling waves that move around the ring with uniform velocity.
5. Synchronization of voltage time curves during bursts goes from  $180^\circ$  anti-phase (low  $g_c$ ), to out-of-phase (intermediate  $g_c$ ), to completely in-phase (high  $g_c$ ). Hence,  $R \sim 1$  is seen only when  $g_c$  is very high ( $> 1000\text{pS}$ ).
6. When multiple cells (of compound and fast or slow) are coupled together, the emergent accordian pattern seen when a pair of cells is coupled, is lost. Irregular bursting and high frequency oscillations are seen (just like the homogeneous case).
7. In a mixture of fast and compound cells, as the percentage of fast cells increases towards 100%, the percentage of time spent bursting by the group of cells tends to that of the case where all the coupled cells are of only fast type.

## Future Directions

1. The effect of heterogeneous coupling has not been studied. It would be interesting to see if and how the results would change by incorporating different coupling strengths among pairs of cells.
2. The 1D ring with periodic boundary conditions can be modified to a 1D ring with fixed boundary conditions.
3. The 1D ring can be extended to 2D or even 3D lattices.

Given that beta cell dysfunction is implicated in Type 2 Diabetes, most likely, models of human beta cells will evolve and develop rapidly. Hopefully, these models will be biologically sufficiently complicated that they will be able to contribute to the design and prediction of new therapies and drugs; while being mathematically sufficiently simple enough that mathematical biologists may analyze and simulate them.

# Chapter 5

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# Appendix A

## MATLAB Programs

The matlab codes (function files and corresponding script files) for the following models have been attached:

- Chay-Keizer single cell model
- Smolen-Keizer single cell model
- Phantom Burster single cell model
- The Dual Oscillator Model

# Code for the Chay-Keizer model

```
1 function dydt = chaykei(t,y)
2     v=y(1);
3     n=y(2);
4     mca=y(3);
5     hca=y(4);
6     ca=y(5);
7
8     % Parameters
9     cm=1;
10    gca=6.5;
11    gk=12;
12    gl=0.04;
13    gkca=0.09; % change to 0.02 to see effects of quinine/TEA
14    vca=100;
15    vk=-75;
16    vl=-40;
17    vprime=50;
18    vstar=30;
19    temp=20;
20    r=8.9e-4;
21    f=0.004;
22    fara=96487;
23    kdiss=1;
24    kca=0.04; %change kca to 0.001 for effects of quinine
25
26    % Gating functions
27    alphamca = -0.1*((v+vprime)-25)/(exp(-(v+vprime)-25)/10)-1);
28    betamca = 4*exp(-(v+vprime)/18);
29    alphahca= 0.07*exp(-(v+vprime)/20);
30    betahca = 1/(exp(-(v+vprime)-30)/10)+1);
31    alphan = (-0.01*((v+vstar)-10)/(exp(-(v+vstar)-10)/10)-1));
32    betan = (0.125*exp(-(v+vstar)/80));
33    phi = 3^((temp-6.3)/10);
34    % To see the effects of TEA use this
35    % n_inf = alphan/(alphan+betan);
36    % tau = 1/(alphan+betan);
37    % taun=tau+2*tau;
38
39    % Ionic currents:
40    i_ca = gca*mca^3*hca*(vca-v);
41    i_k = gk*n^4*(vk-v);
42    i_kca = gkca*ca/(ca+kdiss)*(vk-v);
43    i_l = gl*(vl-v);
44
45    %The differential equations
46    dv = 1/cm*(2*i_ca+i_k+i_kca+i_l);
47    dmca = phi*(alphamca*(1-mca)-betamca*mca);
48    dhca = phi*(alphahca*(1-hca)-betahca*hca);
49    dn = phi*(alphan*(1-n)-betan*n);
50    % Use this dn equation for effect of TEA
51    % dn = phi*(n_inf-n)/taun;
52    dca = f*(3/(r*fara)*i_ca-kca*ca);
53
54    dydt = [dv, dn, dmca, dhca, dca]';
55 end
56 -----
57
```

```

58 %The script file
59
60 %Initial conditions
61
62 v0 = -54.774;
63 n0 = 0.00044035;
64 mca0 = 0.027532;
65 hca0 = 0.086321;
66 ca0 = 0.10749;
67
68 yinit = [v0,n0,mca0,hca0,ca0];
69 tspan = [0 50000];
70 options=odeset('Stats','off');
71
72 tic
73 [t,y] = ode15s(@chaykei,tspan,yinit,options);
74 toc
75
76 figure
77 ax1= subplot(2,1,1)
78 plot(t,y(:,1),'b','LineWidth',2)
79 title('Voltage (V)','FontSize',12)
80 xlabel ('Time (ms)','FontSize',12)
81
82 ax2 = subplot(2,1,2)
83 plot(t,y(:,5),'b','LineWidth',2)
84 title('Ca^{+2}','FontSize',12)
85 xlabel ('Time (ms)','FontSize',12)

```

# Code for the Smolen-Keizer model

```
1 function dydt = smolenkeizer(t,y)
2     v = y(1);
3     n = y(2);
4     ca = y(3);
5     d = y(4);
6     pc = y(5);
7     pot = y(6);
8     ii = y(7);
9     jj = y(8);
10
11     % Parameters
12     vk = -75;
13     captot = 5309;
14     gcam = 1470;
15     x = 0.27;
16     vmj = -11;
17     slopemj = 3.6;
18     gcal = 6.3;
19     cao = 7;
20     gk = 5000;
21     gatpbar = 6000;
22     kone = 0.45;
23     ktwo = 0.012;
24     tot = 1;
25     vn = -20;
26     alambda = 1.2;
27     slopen = 5.3;
28     a = 65;
29     b = 20;
30     c = 60;
31     fr = 0.03;
32     akca = 0.12;
33     rad = 7;
34     farada = 96487;
35     damp = 0;
36     hill = 3;
37     akdamp = 0.6;
38     ka = 2e-05;
39     % change this value to see the effects of glucose
40     r = 0.76;
41     r1 = 0.35;
42     lambdam = 0.78;
43     vm = -15;
44     slopem = 6.2;
45     k2s = 1.512;
46     k2m = 65;
47     lambdak = 1;
48     k3p = 0.02;
49     k3m = 0.005;
50     lambda_i = 1;
51     vi = 36;
52     slope_i = 4.5;
53     tau_i = 2600;
54     lambda_j = 1;
55     vj = 50;
56     slope_j = 6.3;
57     tj = 50000;
```

```

58     tjmin = 1500;
59
60     % activation and time-constant functions
61     phicaj = 1.0/(1.0+exp((vmj-v)/slopemj));
62     phik = 1.0/(1.0+exp((vn-v)/slopen));
63     phica = 1.0/(1.0+exp((vm-v)/slopem));
64     phii = 1.0/(1.0+exp((v+vi)/slopei));
65     phij = 1.0/(1.0+exp((v+vj)/slopej));
66     alpha = 3.0e+03/(8.0e0*pi*farada*rad^3);
67     alphan = exp((v+75.0)/a);
68     betan =exp((-v-75.0)/b);
69     taun = c/(alphan+betan);
70     tauj = tjmin+tj/(exp((v+vj)/(2*slopej))
71             +exp(-(v+vj)/(2*slopej)));
72     gcab = gcam*(x*pot+(1.0-x)*phicaj*jj)+gcal;
73     ghk = cao*v/(1.0-exp(v/13.35));
74     gatp = gatpbar*(1.0+d/kone)/(1.0+d/kone+(tot-d)/ktwo);
75     k1m = lambdam*phica;
76     k1p = lambdam*(1.0-phica);
77     k2p = -k2s*ghk;
78     f2inf = k2p/(k2m+k2p);
79
80     % Ionic currents
81     ica = gcab*ghk;
82     ikv = gk*ii*n*(v-vk);
83     ikatp = gatp*(v-vk);
84
85     % The differential equations
86     dv = -(ica+ikv+ikatp)/captot;
87     dn = alambda*(phik-n)/taun;
88     dca = fr*(-alpha*ica-akca*ca-damp*ca^hill
89             /(akdamp^hill+ca^hill));
90     dd = ka*((tot-d)-d*exp(r*(1.0-(ca/r1))));
91     dpc = k1p*pot*(1.0-f2inf)-k1m*pc;
92     dpot = -(k1p*pot*(1.0-f2inf)-k1m*pc)-
93             lambdak*(k3p*f2inf*pot-k3m*(1.0-pot-pc));
94     dii = lambdai*(phii-ii)/taui;
95     djj = lambdaaj*(phij-jj)/tauaj;
96     dydt =[dv, dn, dca, dd, dpc, dpot, dii, djj]';
97 end
98 -----
99 %The script file
100
101 % Initial conditions
102 v0=-53;
103 n0=0.002;
104 ca0=0.15;
105 d0=0.42;
106 pc0=0.9;
107 pot0=0.02;
108 ii0=0.98;
109 jj0=0.47;
110
111 yinit = [v0,n0,ca0,d0,pc0,pot0,ii0,jj0];
112 tspan = [0 50000];
113 options=odeset('Stats','off');
114

```



```
115 tic
116 [t,y] = ode15s(@smolenkeizer,tspan,yinit,options);
117 toc
118
119 figure
120 ax1 = subplot(2,1,1)
121 plot(t,y(:,1),'b','LineWidth',2)
122 title('Voltage (mV)','FontSize',20);
123 set(gca,'FontSize',17) % xlabel('Time (ms)');
124
125 ax2 = subplot(2,1,2)
126 plot(t,y(:,3),'b','LineWidth',2)
127 title('[Ca+2'],'FontSize',20);
128 xlabel('Time (ms)','FontSize',20);
129 set(gca,'FontSize',17)
```

# Code for the Phantom Burster model

```
1 % Phantom bursting model, with 2 fast and 2 slow variables.
2 % Units: V = mV; t = ms; g = pS; I = fA
3
4 function dydt = phantomburster(t,y)
5     v = y(1);
6     n = y(2);
7     s1 = y(3);
8     s2 = y(4);
9
10    % Parameters
11    vl = -40;
12    vca = 100;
13    vk = -80;
14    lambda = 1.1;
15    gca = 280;
16    gk = 1300;
17    gl = 25;
18    vs1 = -40;
19    taus1 = 1000; %could go upto 10s
20    vs2 = -42;
21    taus2 = 120000; %could be as small as 1min
22    gs2 = 32;
23    gs1 = 3; % for slow bursting
24    % gs1 = 4; %for medium bursting 'phantom bursting'
25    % gs1 = 20; % for fast bursting
26    cm = 4524;
27    tnbar = 9.09;
28    vm = -22;
29    vn = -9;
30    sm = 7.5;
31    sn = 10;
32    ss1 = 0.5;
33    ss2 = 0.4;
34
35    % activation and time-constant functions
36    minf = 1.0/(1.0+exp((vm-v)/sm));
37    ninf = 1.0/(1.0+exp((vn-v)/sn));
38    taun = tnbar/(1.0+exp((v-vn)/sn));
39    slinf = 1.0/(1.0+exp((vs1-v)/ss1));
40    s2inf = 1.0/(1.0+exp((vs2-v)/ss2));
41
42    % Ionic currents
43    ica = gca*minf*(v-vca);
44    ik = gk*n*(v-vk);
45    il = gl*(v-vl);
46    is1 = gs1*s1*(v-vk);
47    is2 = gs2*s2*(v-vk);
48
49    % The differential equations
50
51    if 15000<t && t<45000
52        istim = 0;
53    else
54        istim=0;
55    end
56
57    dv = -( ica + ik + il + is1 + is2 )/cm - istim/cm;
```

```

58     dn = lambda*(ninf - n)/taun;
59     ds1 = (s1inf - s1)/taus1;
60     ds2 = (s2inf - s2)/taus2;
61
62     dydt = [dv, dn, ds1, ds2]';
63 end
64 -----
65 % The script file
66
67 % Initial conditions
68
69 v0 = -43.0;
70 n0 = 0.03;
71 s10 = 0.1;
72 s20 = 0.434;
73
74 yinit = [v0,n0,s10,s20];
75 tspan = [0 100000];
76 options=odeset('Stats','off');
77
78 tic
79 [t,y] = ode15s(@phantomburster,tspan,yinit,options);
80 toc
81
82 figure(1);
83 % t(t<3e4) = NaN;
84 plot(t,y(:,1),'b','LineWidth',2)
85 title('Fast bursting - Voltage (mV)', 'FontSize',21)
86 xlabel('Time (ms)', 'FontSize',21);
87 set(gca, 'FontSize',25)

```

# Code for the Dual Oscillator Model

```
1 % Program for a beta-cell model coupled to glycolysis -DOM
2 % This was published by Bertram, Sherman et. al.
3 % Biophysical Journal, 87:3074-3087, Nov. 2004.
4
5 % State variables: % V    -- membrane potential
6 % n    -- activation of delayed rectifier
7 % Ca   -- free cytosolic calcium concentration
8 % ADP  -- cytosolic ADP concentration
9 % CaER -- concentration of free calcium in the ER
10 % G6P  -- glucose 6-phosphate concentration
11 % FBP  -- fructose 1,6-bisphosphate concentration
12
13 function dydt = dualoscillator(t,y)
14
15     % y = [V, n, Ca, CaER, ADP, G6P, FBP];
16     V = y(1);
17     n = y(2);
18     Ca = y(3);
19     CaER = y(4);
20     ADP = y(5);
21     G6P = y(6);
22     FBP = y(7);
23
24     % Parameter sets for various behaviors:
25     % k_gamma = 10 for all of the sets!
26     % Note: An inc. in R_GK is like an inc in gluc.
27
28     % for compound bursting
29     % R_GK = 0.2; gKATPbar = 25000; gKCabar = 600;
30
31     % for slow bursting
32     % R_GK = 0.2; gKATPbar = 27000; gKCabar = 100;
33
34     % for subthreshold bursting
35     % R_GK = 0.2; gKATPbar = 30000; gKCabar = 100;
36
37     % for accordian bursting
38     %R_GK = 0.2; gKATPbar = 23000; gKCabar = 600;
39
40     % for fast bursting
41     % R_GK = 0.4; gKATPbar = 25000; gKCabar = 600;
42
43     % -----
44     % Channel properties
45
46     Cm = 5300;
47     VK = -75;
48     gK = 2700;
49     taun = 20;
50     ninf = 1/(1+exp(-(16+V)/5));
51     gCa = 1000;
52     VCa = 25;
53     minf = 1/(1+exp(-(20+V)/12));
54     % gKCabar = 600; (uncomment out one of the sets!)
55     Kd = 0.5;
56
57     % Ionic currents:
```

```

58 IK = gK*n*(V-VK);
59 ICa = gCa*minf*(V-VCa);
60 gKCa = gKCabar/(1+(Kd/Ca)^2);
61 IKCa = gKCa*(V-VK);
62
63 % IKATP (see below, after calc. of nucleotide concs.)
64 % gKATPbar = 25000;(uncomment out one of the sets!)
65
66 % Calcium Handling
67 alpha = 4.50e-6;
68 kPMCA = 0.2;
69 fcyt = 0.01;
70 fer = 0.01;
71 % sigmav=cyt volume/ER volume = V_cyt/V_er
72 sigmav = 31; pleak = 0.0002;
73 kSERCA = 0.4;
74 Jmem = -(alpha*ICa + kPMCA*Ca);
75 JSERCA = kSERCA*Ca;
76 Jleak = pleak*(CaER - Ca);
77 Jer = Jleak - JSERCA;
78
79 % -----
80 % Glycolytic and Keizer-Magnus components
81 % Parameters
82 % R_GK--glucokinase rate
83 % Atot--total adenine nucleotide concentration (micromolar)
84 % K1--Kd for AMP binding
85 % K2--Kd for FBP binding
86 % K3--Kd for F6P binding
87 % K4--Kd for ATP binding
88 % famp, etc--Kd amplification factors for heterotropic binding
89 % famp corresponds to f13 in the paper
90 % fmt corresponds to f41 in the paper
91 % ffbp corresponds to f23 in the paper
92 % fbt corresponds to f42 in the paper
93 % fatp corresponds to f43 in the paper
94 % R_GPDH--glyceraldehyde phosphate dehydrogenase rate
95
96 % Glycolytic parameters
97 K1 = 30;
98 K2 = 1;
99 K3 = 50000;
100 K4 = 1000;
101 famp = 0.02;
102 fmt = 20;
103 ffbp = 0.2;
104 fbt = 20;
105 fatp = 20;
106
107 % Glycolytic expressions
108 F6P = 0.3*G6P;
109 % nucleotide concentrations used for R_PFK
110 Atot = 3000;
111 rad = sqrt((ADP-Atot)^2-4*ADP^2);
112 ATP = 0.5*(Atot-ADP+rad);
113 % Use this to see how the glyc. comp. osc. independently!
114 % ATP = 2000;

```

```

115 AMP = ADP^2/ATP;
116
117 % Iterative calculation of R_PFK
118 % (cf. Smolen95, Eq. 12)
119 % alpha = 1 -- AMP bound
120 % beta = 1 -- FBP bound
121 % gamma = 1 -- F6P bound
122 % Δ = 1 -- ATP bound
123 % (alpha,beta,gamma,Δ)
124
125 % (0,0,0,0)
126 weight1 = 1;
127 topa1 = 0;
128 bottom1 = weight1;
129
130 % (0,0,0,1)
131 weight2 = ATP^2/K4;
132 topa2 = topa1;
133 bottom2 = bottom1+weight2;
134
135 % (0,0,1,0)
136 weight3 = F6P^2/K3;
137 topa3 = topa2+weight3;
138 bottom3 = bottom2+weight3;
139
140 % (0,0,1,1)
141 weight4 = (F6P*ATP)^2/(fatp*K3*K4);
142 topa4 = topa3+weight4;
143 bottom4 = bottom3+weight4;
144
145 % (0,1,0,0)
146 weight5 = FBP/K2;
147 topa5 = topa4;
148 bottom5 = bottom4+weight5;
149
150 % (0,1,0,1)
151 weight6 = (FBP*ATP^2)/(K2*K4*fbt);
152 topa6 = topa5;
153 bottom6 = bottom5+weight6;
154
155 % (0,1,1,0)
156 weight7 = (FBP*F6P^2)/(K2*K3*ffbp);
157 topa7 = topa6+weight7;
158 bottom7 = bottom6+weight7;
159
160 % (0,1,1,1)
161 weight8 = (FBP*F6P^2*ATP^2)/(K2*K3*K4*ffbp*fbt*fatp);
162 topa8 = topa7+weight8;
163 bottom8 = bottom7+weight8;
164
165 % (1,0,0,0)
166 weight9 = AMP/K1;
167 topa9 = topa8;
168 bottom9 = bottom8+weight9;
169
170 % (1,0,0,1)
171 weight10 = (AMP*ATP^2)/(K1*K4*fmt);

```

```

172     topa10 = topa9;
173     bottom10 = bottom9+weight10;
174
175     % (1,0,1,0)
176     weight11 = (AMP*F6P^2)/(K1*K3*famp);
177     topa11 = topa10+weight11;
178     bottom11 = bottom10+weight11;
179
180     % (1,0,1,1)
181     weight12 = (AMP*F6P^2*ATP^2)/(K1*K3*K4*famp*fmt*fatp);
182     topa12 = topa11+weight12;
183     bottom12 = bottom11+weight12;
184
185     % (1,1,0,0)
186     weight13 = (AMP*FBP)/(K1*K2);
187     topa13 = topa12;
188     bottom13 = bottom12+weight13;
189
190     % (1,1,0,1)
191     weight14 = (AMP*FBP*ATP^2)/(K1*K2*K4*fbt*fmt);
192     topa14 = topa13;
193     bottom14 = bottom13+weight14;
194
195     % (1,1,1,0) -- the most active state of the enzyme
196     weight15 = (AMP*FBP*F6P^2)/(K1*K2*K3*ffbp*famp);
197     topa15 = topa14;
198     topb = weight15;
199     bottom15 = bottom14+weight15;
200
201     % (1,1,1,1)
202     weight16 = ...
           (AMP*FBP*F6P^2*ATP^2)/(K1*K2*K3*K4*ffbp*famp*fbt*fmt*fatp);
203     topa16 = topa15+weight16;
204     bottom16 = bottom15+weight16;
205
206     % Phosphofructokinase rate
207     % lambda, Vmax as in Smolen95, Eq. 3
208     lambda = 0.06; %in the paper lambda <<1
209     Vmax = 2;
210     R_PFK = Vmax*(lambda*topa16 + topb)/bottom16;
211     % GPDH flux:
212     R_GPDH = 0.2*sqrt(FBP);
213
214     % KATP channel
215     Kdd = 17;
216     Ktd = 26;
217     Ktt = 1;
218
219     % KATP channel open probability
220     %(reference: Magnus and Keizer (1998))
221     mgADP = 0.165*ADP;
222     ADP3m = 0.135*ADP;
223     ATP4m = 0.05*ATP;
224     topo = 0.08*(1+2*mgADP/Kdd) + 0.89*(mgADP/Kdd)^2;
225     bottomo = (1+mgADP/Kdd)^2 * (1+ADP3m/Ktd+ATP4m/Ktt);
226     oinf = topo/bottomo;    gKATP = gKATPbar*oinf;
227     IKATP = gKATP*(V-VK);

```

```

228
229     % glycolytic input to mitochondrial ADP equation:
230     k_gamma = 10;
231     v_gamma = 2.2;
232     gamma = v_gamma*(R_GPDH/(k_gamma+R_GPDH));
233     r = 1;
234     taua = 300000;
235     r1 = 0.35;
236
237     % conversion parameter for glycolytic subsystem
238     % kappa erroneously called lambda in paper;
239     % renamed kappa for consistency with Smolen,
240     %JTB, 1995 and Pedersen et al, 2005.
241     kappa = 0.005;
242     % R_GK = 0.2; (uncomment out one of the sets!)
243
244     % Differential equations
245
246     dV = -(IK + ICa + IKCa + IKATP)/Cm;
247     dn = (ninf-n)/taun;
248     dCa = fcyt*(Jmem + Jer);
249     dCaER = -fer*sigmav*Jer;
250     dADP = (ATP-ADP*exp((r + gamma)*(1-Ca/r1)))/taua;
251     dG6P = kappa*(R_GK - R_PFK);
252     dFBP = kappa*(R_PFK - 0.5*R_GPDH);
253
254     dydt = [dV, dn, dCa, dCaER, dADP, dG6P, dFBP]';
255 end
256 -----
257 %The script file
258
259 % Initial conditions:
260
261 V0 = -60;
262 n0 = 0;
263 Ca0 = 0.1;
264 CaER0 = 185;
265 ADP0 = 780;
266 G6P0 = 200;
267 FBP0 = 40;
268
269 yinit = [V0,n0,Ca0,CaER0,ADP0,G6P0,FBP0];
270 tspan = [0 1000000];
271 options=odeset('Stats','off');
272
273 tic
274 [t,y] = ode15s(@dualoscillator,tspan,yinit,options);
275 toc
276
277 % figure
278 % plot(t,y(:,1),'b','LineWidth',2)

```